

November 9-15, 2019

An-Najah N. University, Nablus, Palestine

Detection of gravitational waves

- Gravitational-wave detectors
- LIGO-Virgo detector network
- Signal extraction methods
- Detector's noise



LIGO Hanford, USA



Virgo Pisa, Italy



LIGO Livingston, USA

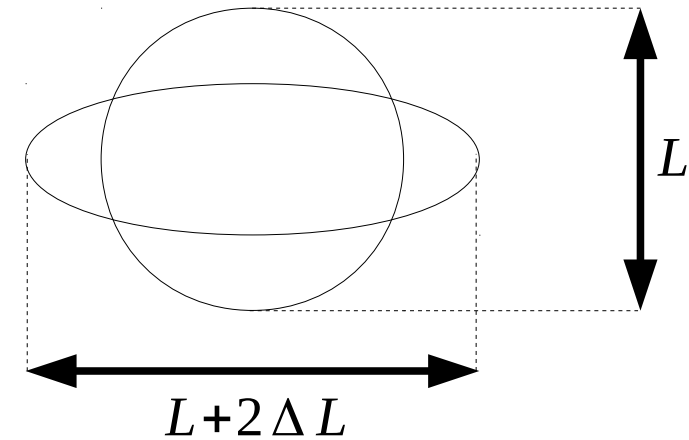
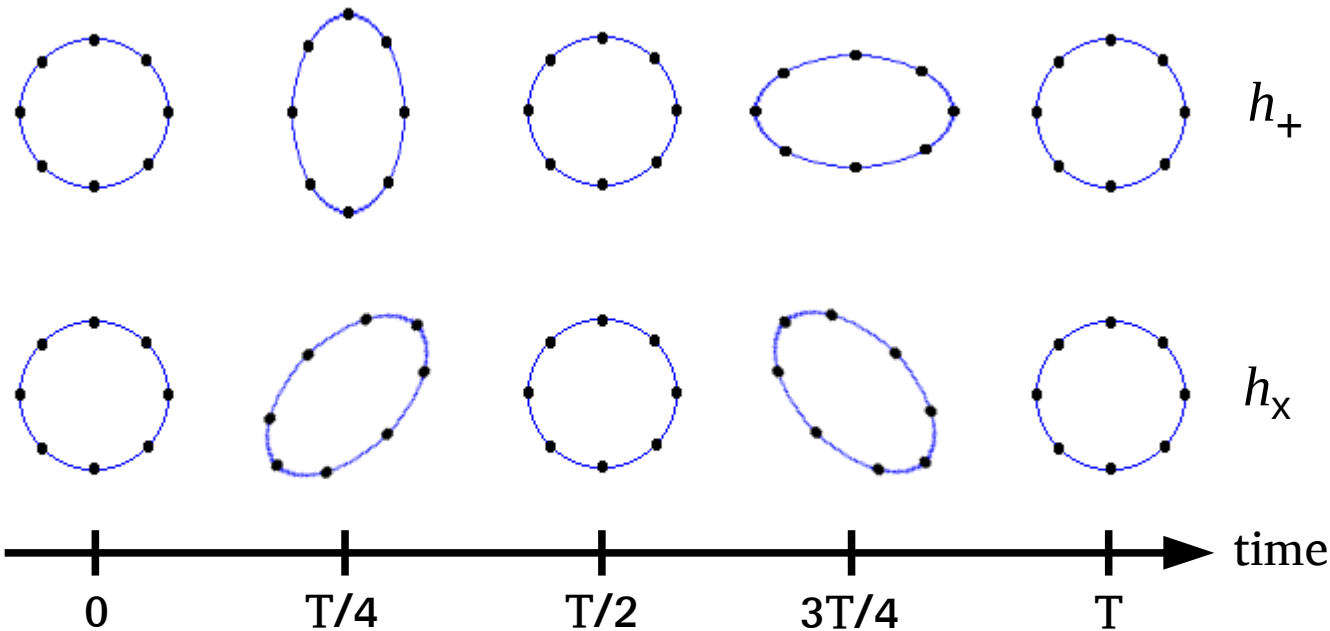


Gravitational wave detection

Add a small perturbation to the Minkowski metric: $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$ $|h_{\mu\nu}| \ll 1$

- h obeys a plane-wave equation
- the wave propagates at the speed of light
- 2 degrees of freedom: h_+ and h_\times

→ Gravitational waves



$$h = 2 \frac{\Delta L}{L}$$

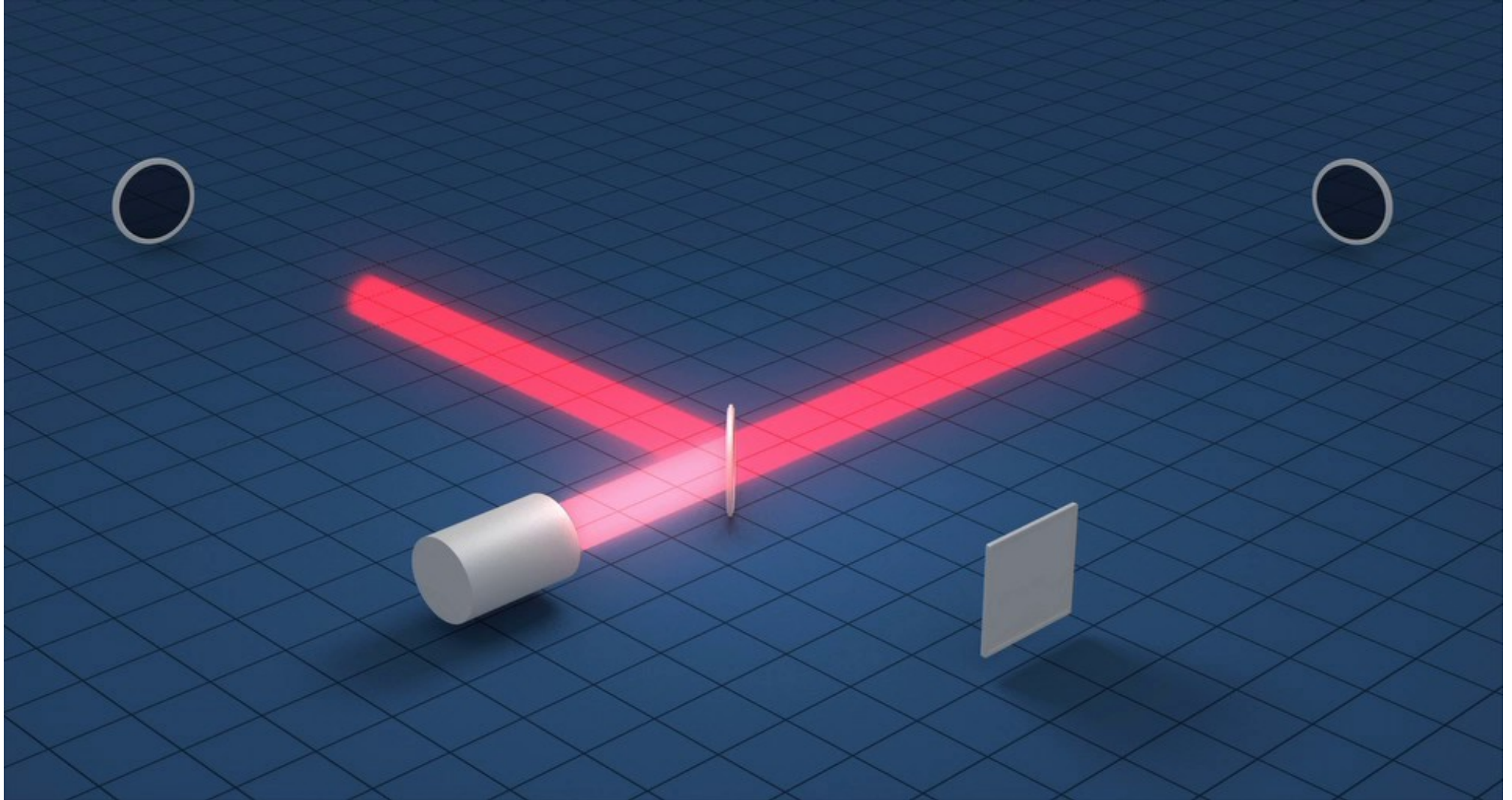
When they propagate, gravitational waves

- do not interact with matter
- are attenuated by $1/r$

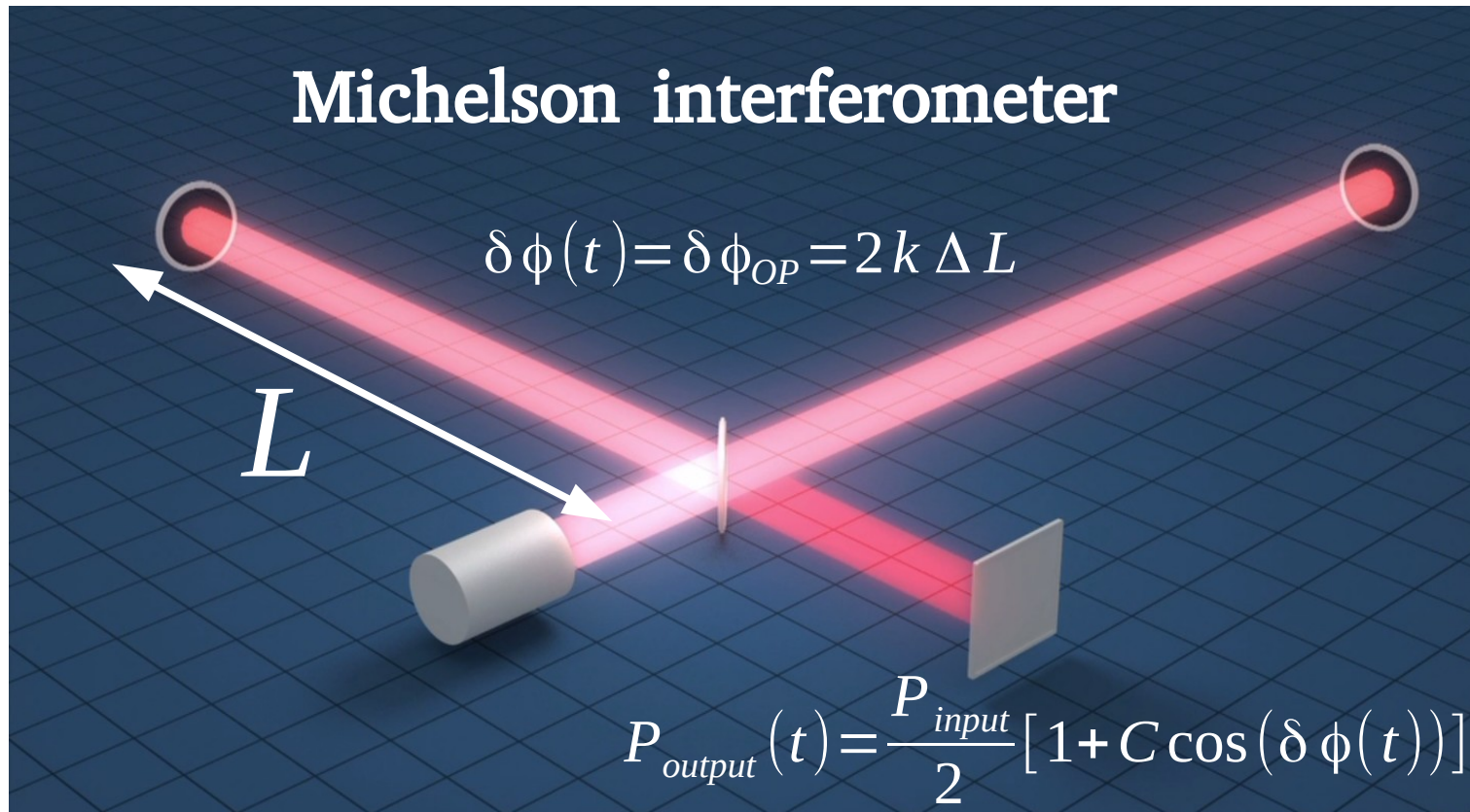
→ **Gravitational waves are the perfect probe!**

→ **BUT... $h \sim 10^{-21}$**

Gravitational wave detectors

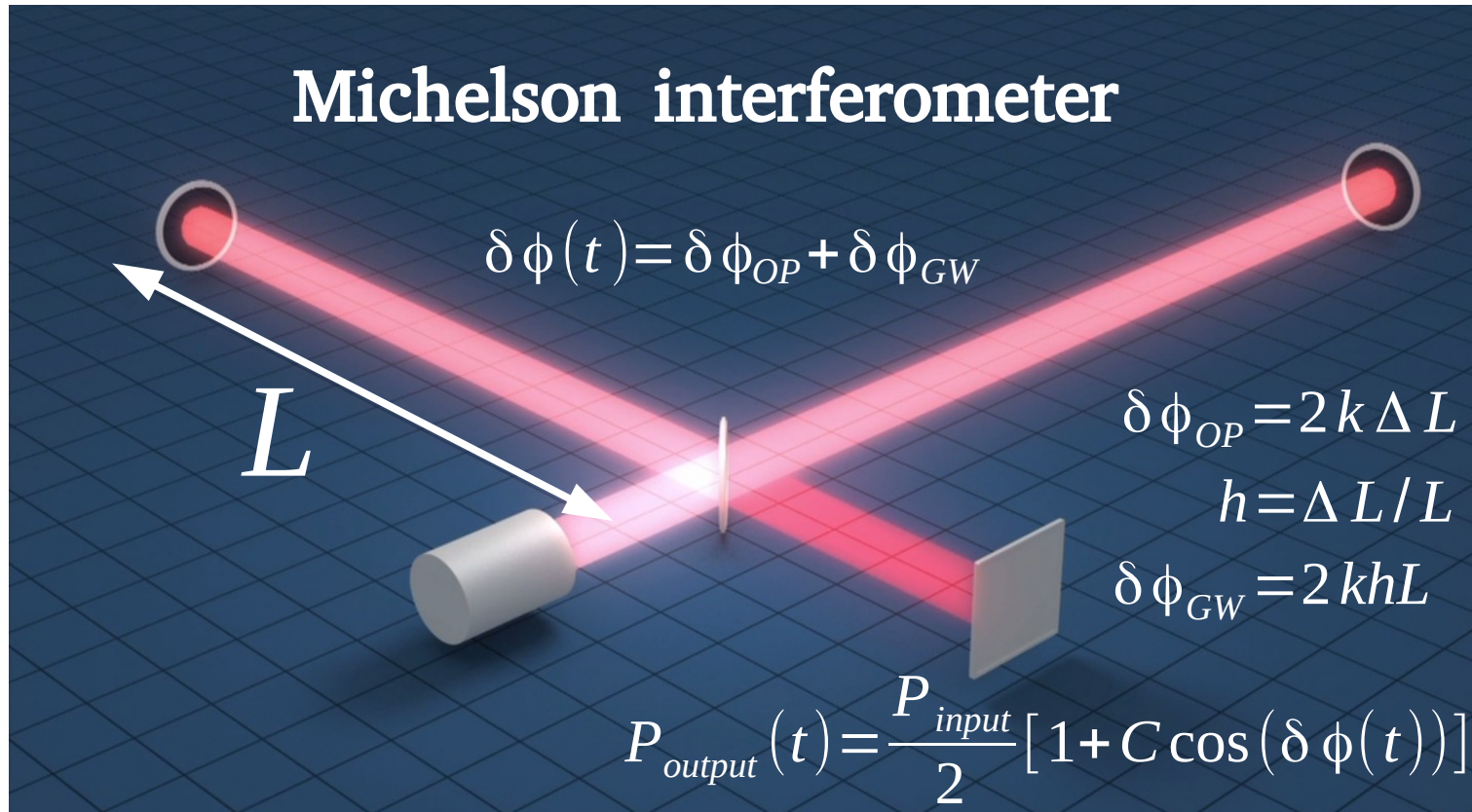


Gravitational wave detectors



Operation: set output on a dark fringe

Gravitational wave detectors



$$P_{output}(t) \simeq \frac{P_{input}}{2} [1 + C \cos(\delta \phi_{OP}) - C \sin(\delta \phi_{OP}) \times \delta \phi_{GW}(t)]$$

→ A gravitational wave is detected as a power variation

Sensitivity limited by noise

The detector's sensitivity to h is limited by noise

For example: shot noise due to uncertainty in photon counting rate

$$N_{\text{photon}} \propto P_{\text{output}} \quad \rightarrow \quad \delta N_{\text{photon}} \propto \sqrt{N_{\text{photon}}} \quad \rightarrow \quad \delta P_{\text{shotnoise}} \propto \sqrt{P_{\text{output}}}$$

Signal-to-noise ratio

$$\frac{S}{B} = \frac{\delta P_{\text{GW}}}{\delta P_{\text{shot noise}}} \propto L \sqrt{P_{\text{input}}} h$$

$$\text{If } \frac{S}{B} = 1 \quad \rightarrow \quad h_{\text{shot noise}} \propto \frac{1}{L \sqrt{P_{\text{input}}}}$$

Table-top Michelson interferometer $\rightarrow 10^{-17}$

Astrophysical sources $\rightarrow 10^{-21}$

Boosting the sensitivity

$$\delta P_o = G_{PR} P_i C \frac{2\pi}{\lambda} \sin\left(\frac{4\pi}{\lambda} \Delta L_0\right) \frac{2F}{\pi} \delta \Delta L$$

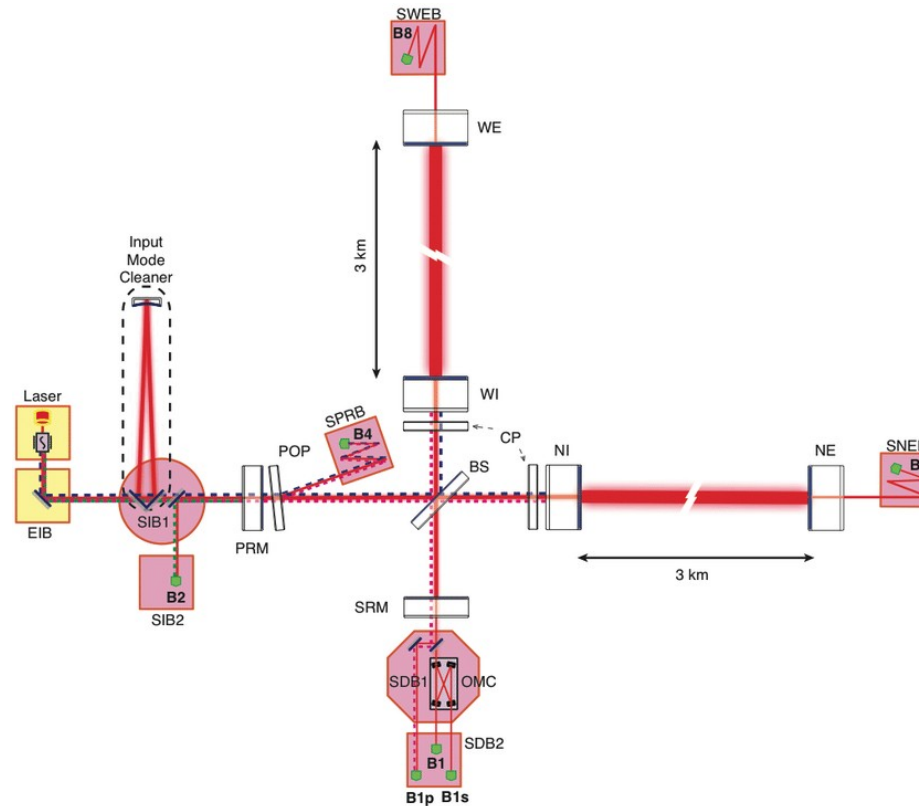
Typical shot noise at 50 mW:

$$\delta P_{shot} \sim 0.1 \text{ nW}$$

$$\delta \Delta L \sim 5 \times 10^{-20} \text{ m}$$

Strain amplitude (reconstruction)

$$h = \frac{\delta \Delta L}{L_0} \sim 10^{-23}$$



$$\lambda = 1064 \text{ nm}$$

$$P_i = 100 \text{ W}$$

$$\Delta L_0 = 10^{-11} \text{ m}$$

$$F = 450$$

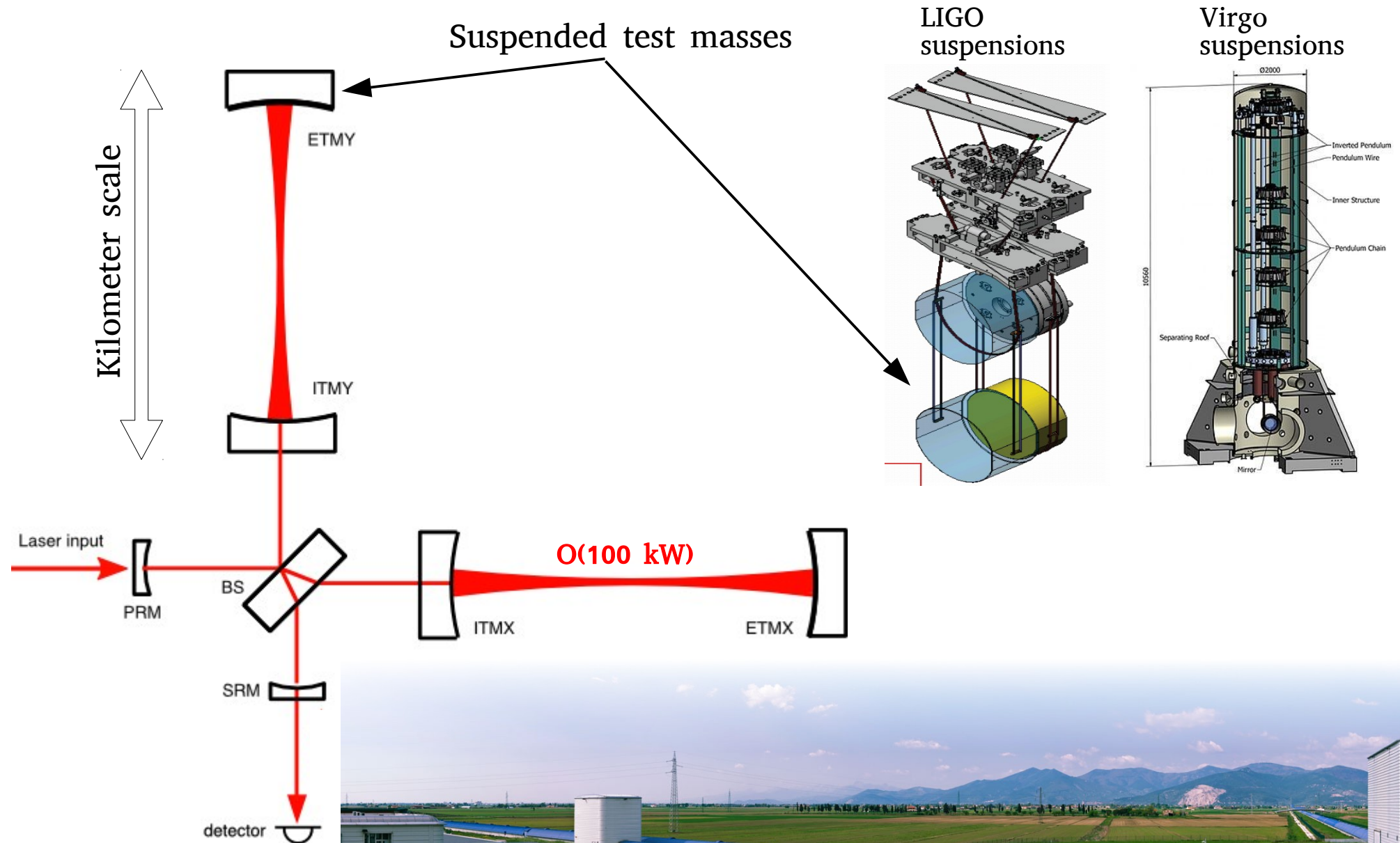
$$C = 1$$

$$G_{PR} = 38$$



In reality the reconstruction is frequency dependent

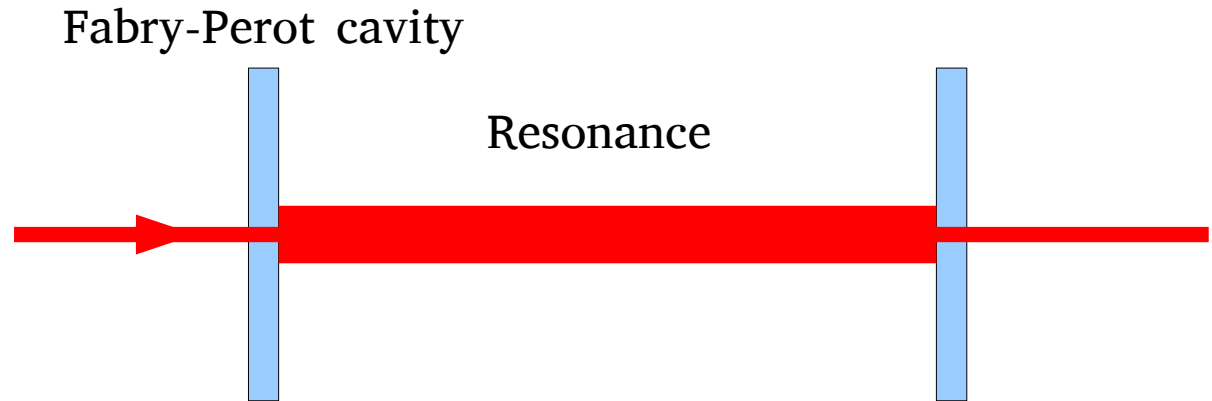
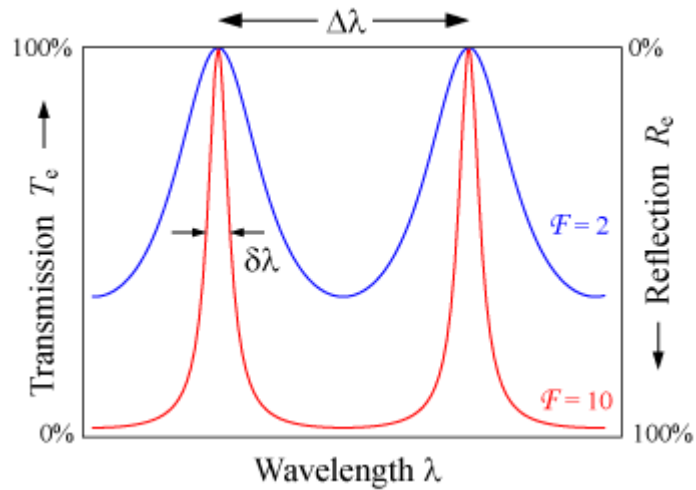
Gravitational wave detectors



Virgo interferometer

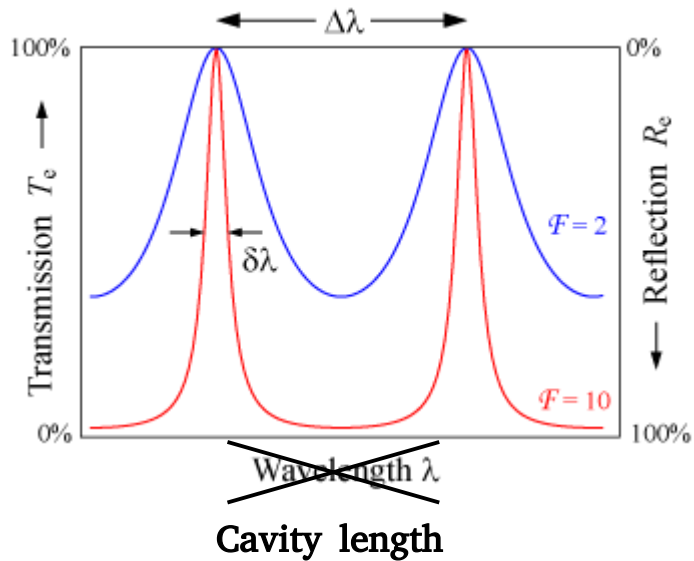
Lock control

The detector's mirrors must be "controlled" to lock and maintain the cavities at resonance

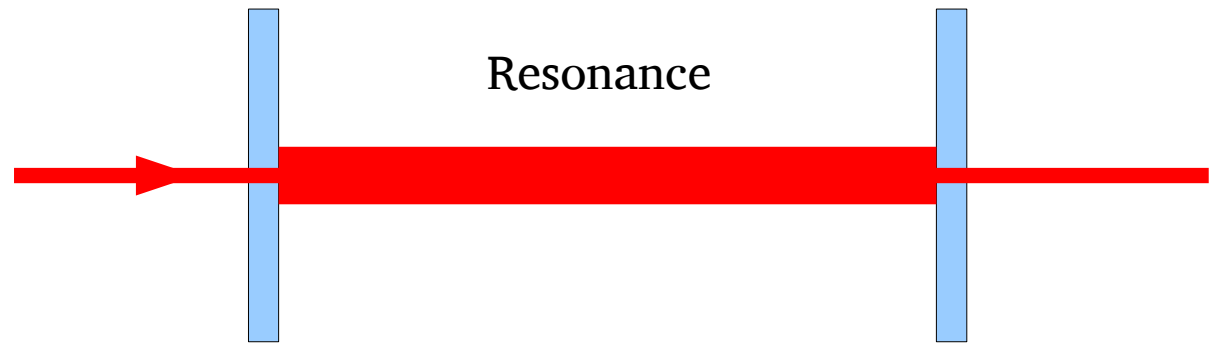


Lock control

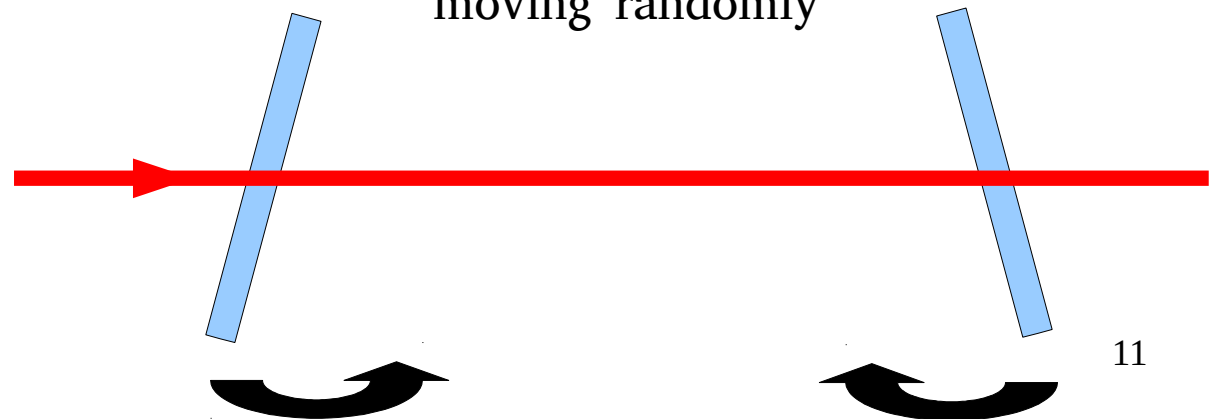
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Fabry-Perot cavity

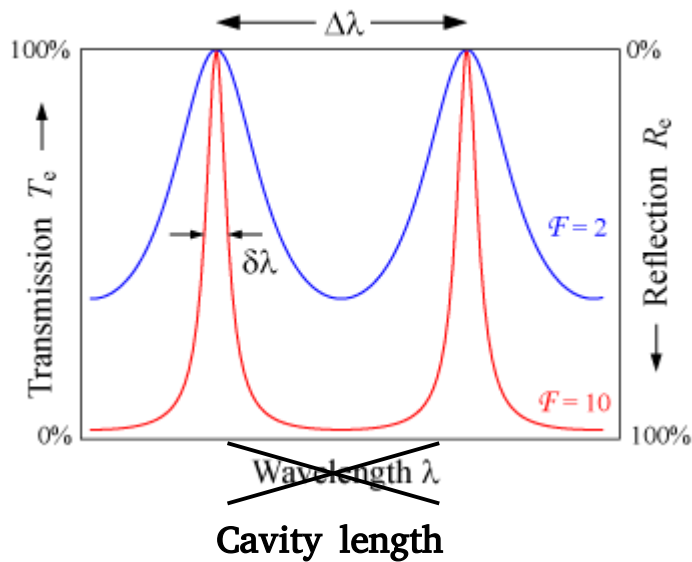


Suspended mirrors moving randomly



Lock control

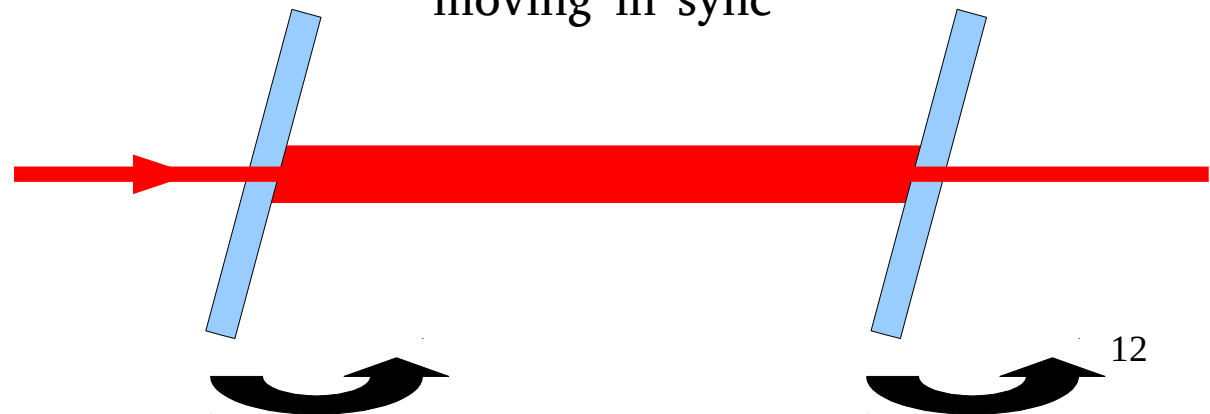
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Fabry-Perot cavity



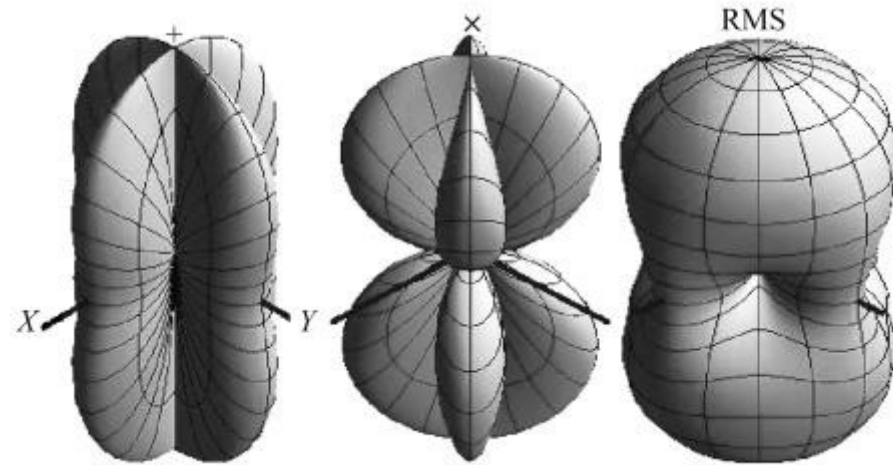
Suspended mirrors moving in sync



A control loop is activated to bring (and maintain) the mirrors at resonance

Sky coverage

The detector's sensitivity over the sky is not uniform



$$h_{det}(t) = F_{+}(t, \underline{ra}, \underline{dec}, \underline{\Psi}) \times h_{+}(t) + F_{\times}(t, \underline{ra}, \underline{dec}, \underline{\Psi}) \times h_{\times}(t)$$

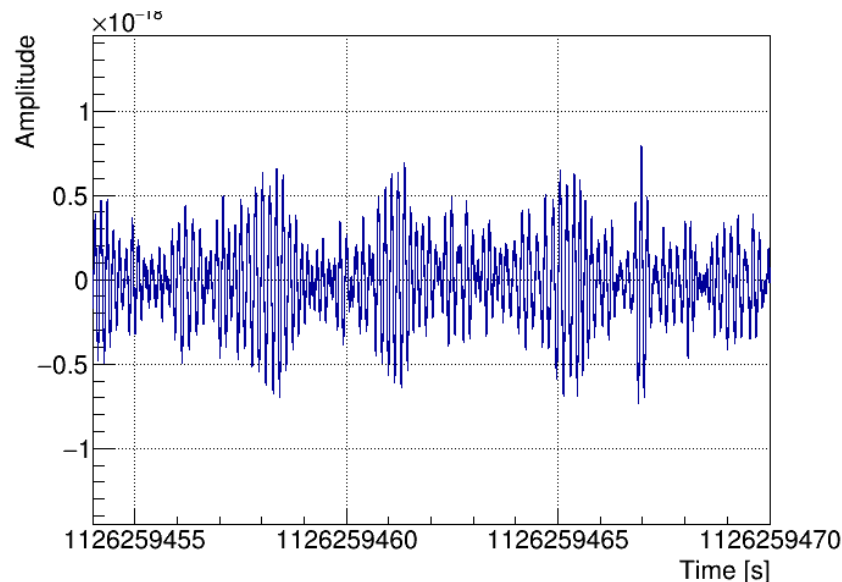
Source position Source polarization angle

Gravitational-wave data

GW detectors' readout system provides at any instant an estimate of strain: a quantity that is sensitive to arms' length difference:

→ Digitized discrete time series: $raw(t)$ (sampled at 16384 Hz or 20000 Hz) and synchronized with GPS clocks.

→ Calibration of $raw(t)$: apply a frequency dependent factor [in reality this is a bit more complicated ...]



→ $h_{det}(t)$ time series that is detector noise plus all hypothetical GW signals

$$h_{det}(t) = n(t) + \mathbf{GW}(t)$$

Power spectral density

Fourier transform

A time series $s(t)$ can be projected over a basis of sinusoidal functions:

$$\tilde{s}(f) = \int_{-\infty}^{\infty} s(t) e^{-2i\pi ft} dt \quad (\text{forward})$$

$$s(t) = \int_{-\infty}^{\infty} \tilde{s}(f) e^{2i\pi ft} df \quad (\text{backward})$$

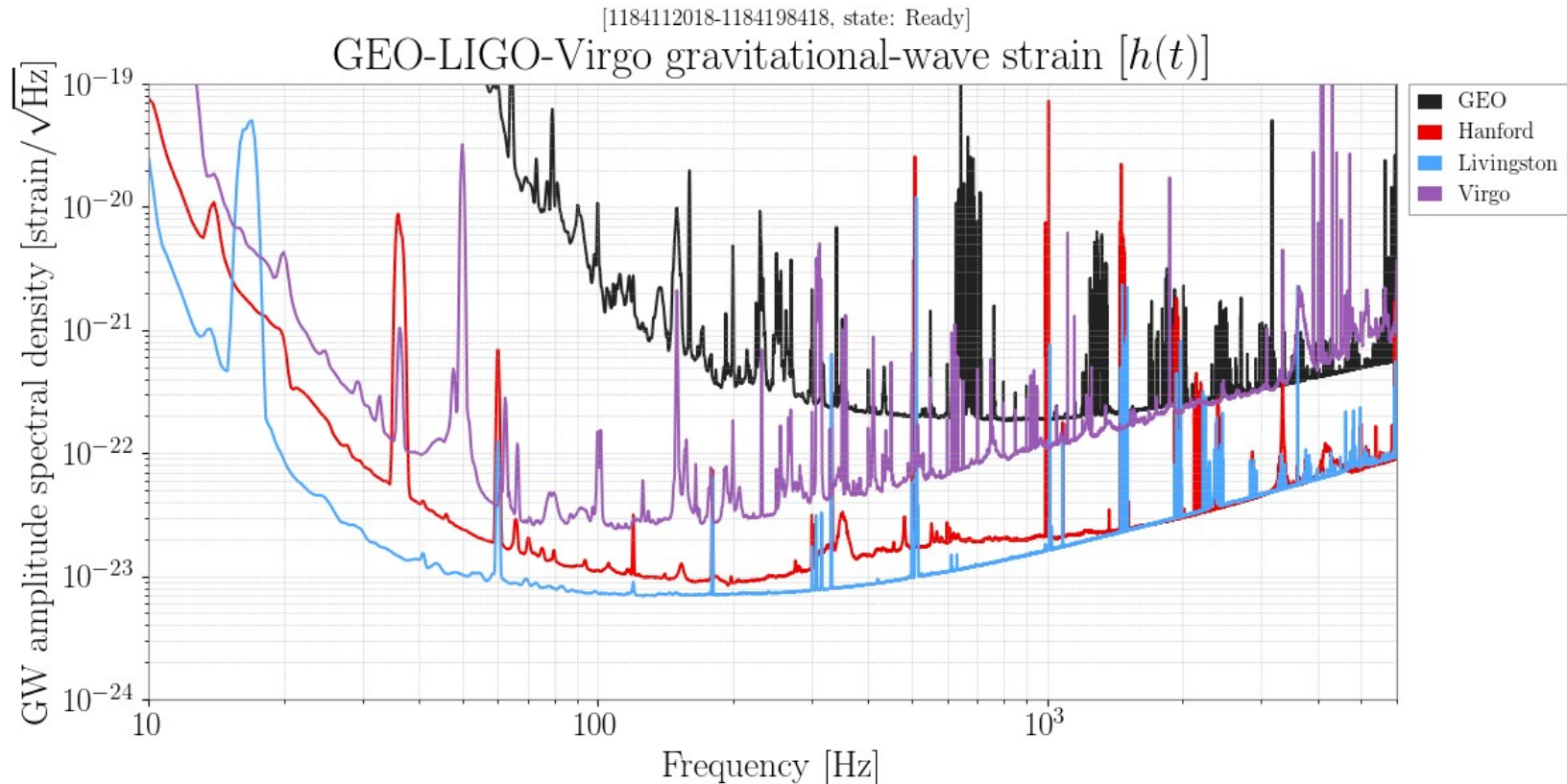
The signal is decomposed in characteristic frequencies

A noise source $n(t)$ limiting the extraction of a signal $s(t)$ is completely characterized by the power (amplitude) spectral density $S_n(f)$

$$S_n(f) = 2|\tilde{n}(f)|^2 \quad A_n(f) = \sqrt{S_n(f)}$$

Detector sensitivity

LIGO-Virgo sensitivity – 2017



Fundamental noises:

- $f < 10$ Hz → seismic noise
- 10 Hz $< f < 200$ Hz → thermal noise
- $f > 200$ Hz → quantum shot noise

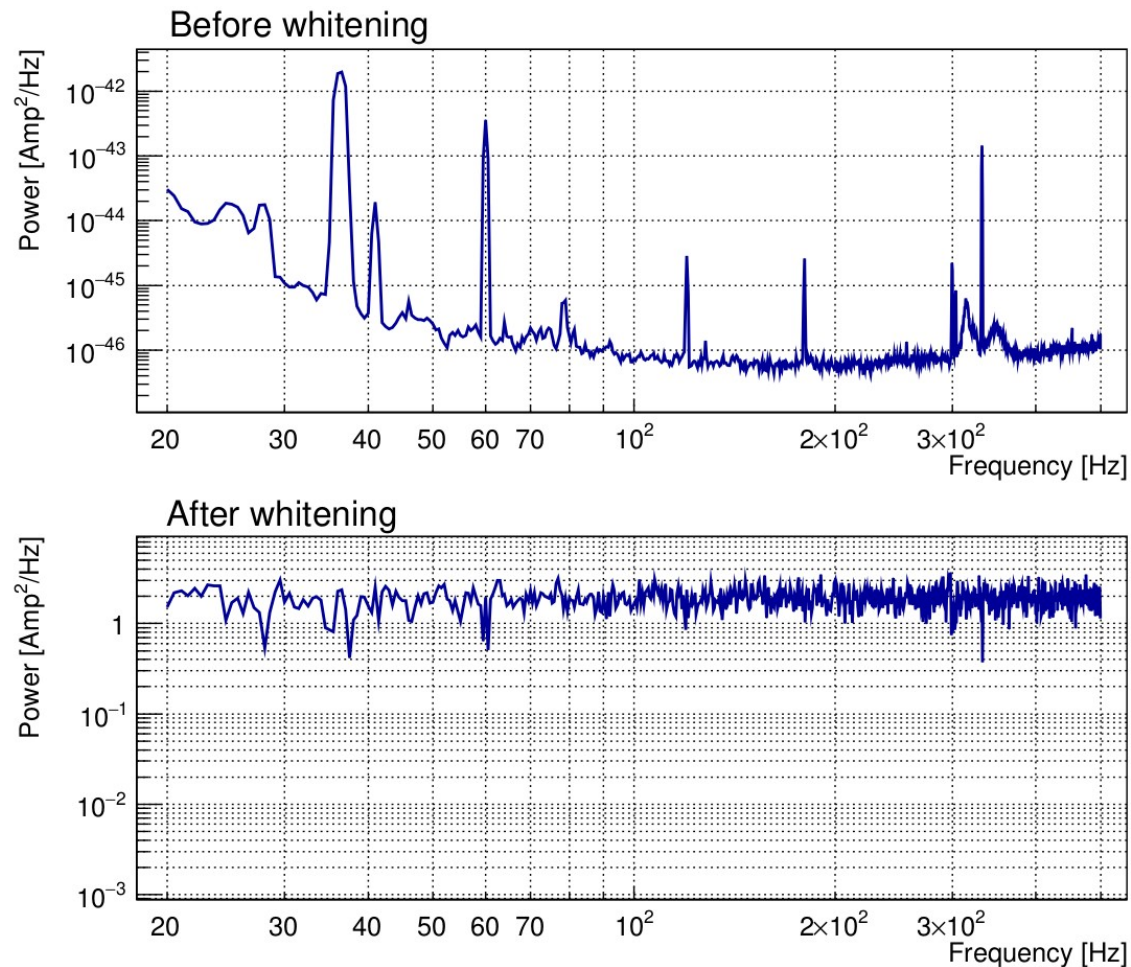
+ technical noise (environment, scattered light, control...)

Data whitening

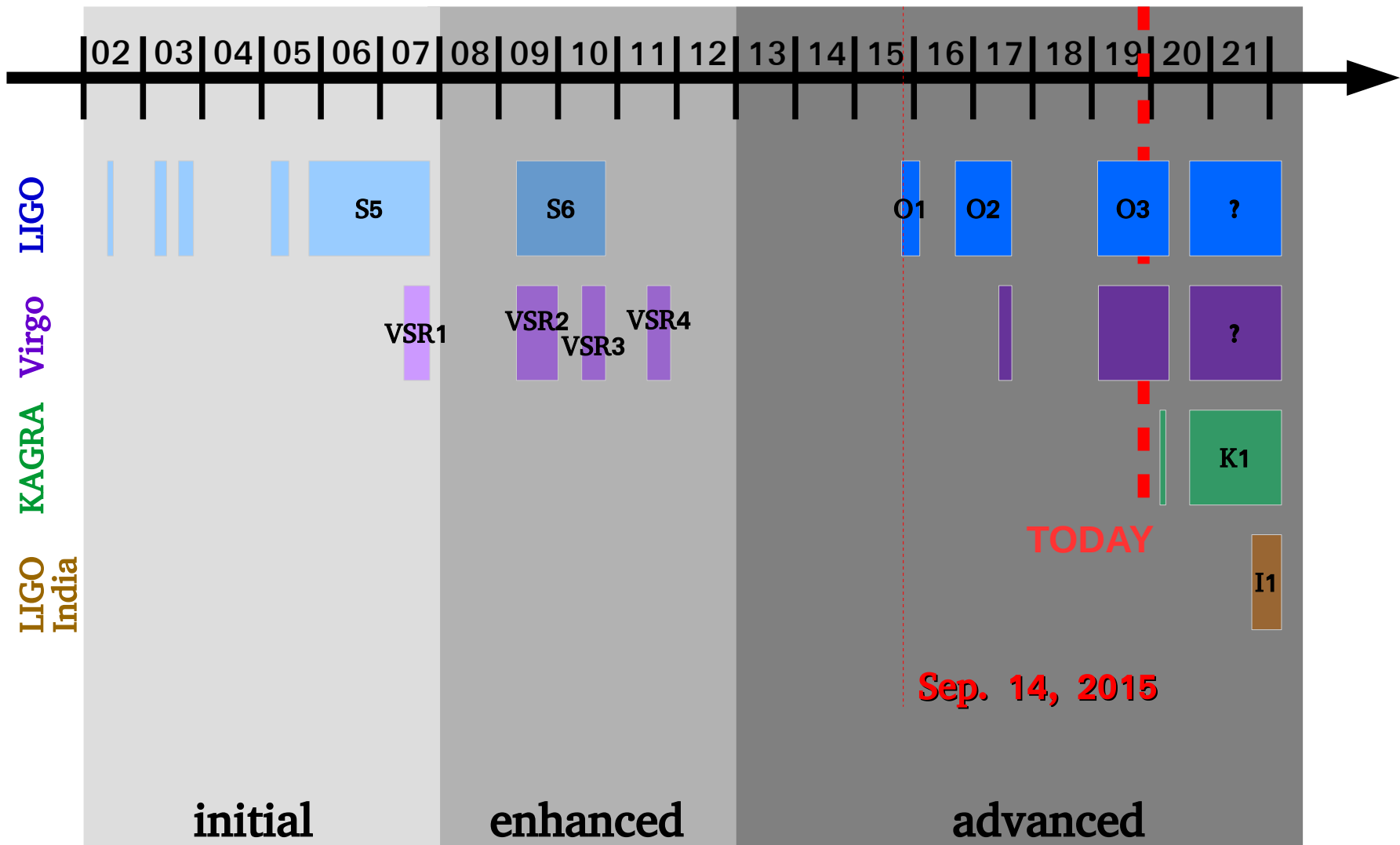
GW data must be whitened. Several methods are used :

- reweighting of frequency bins
- linear prediction

→ white noise is mandatory for statistical interpretation of the data



Scientific runs



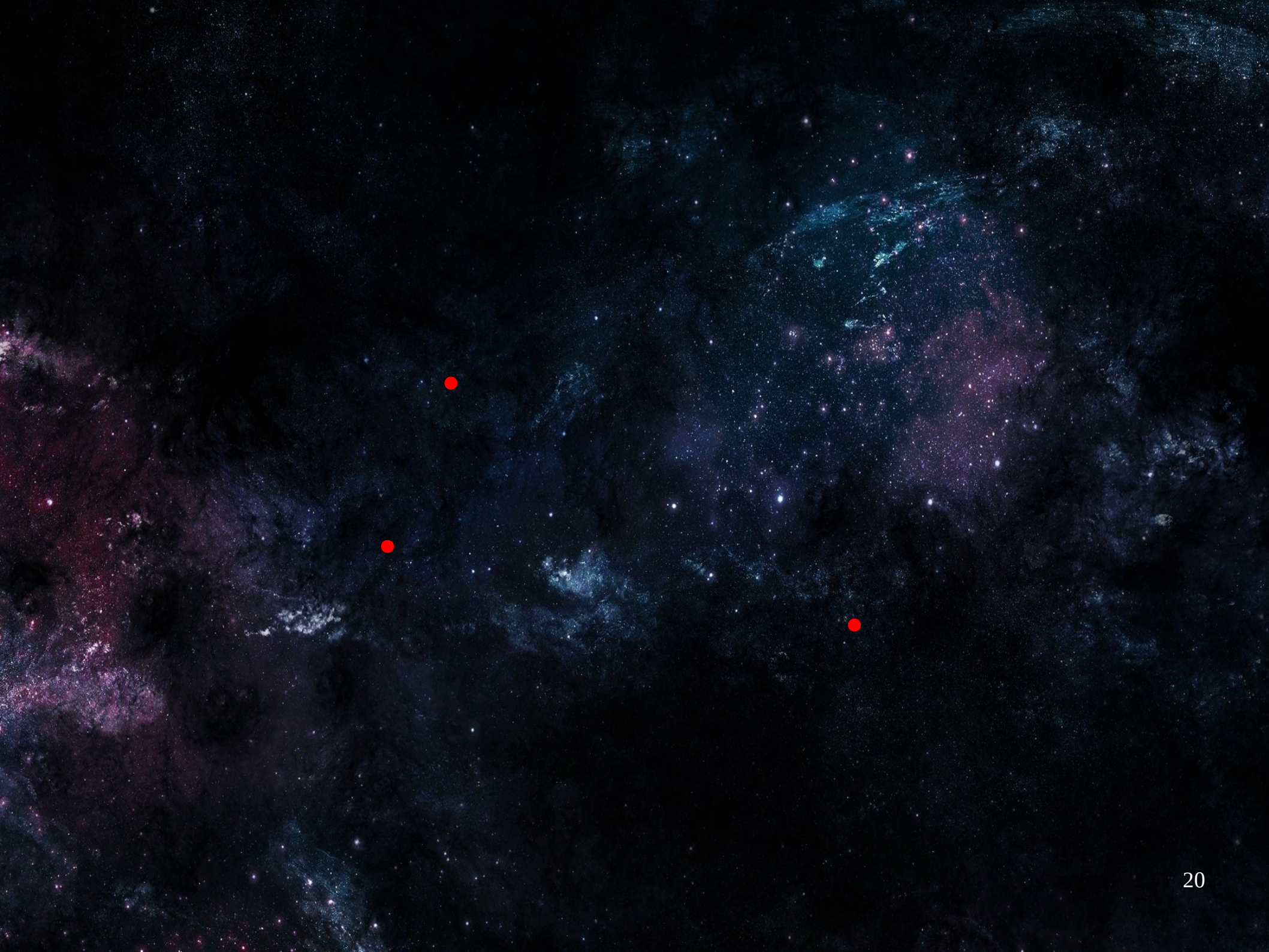
Data analysis :

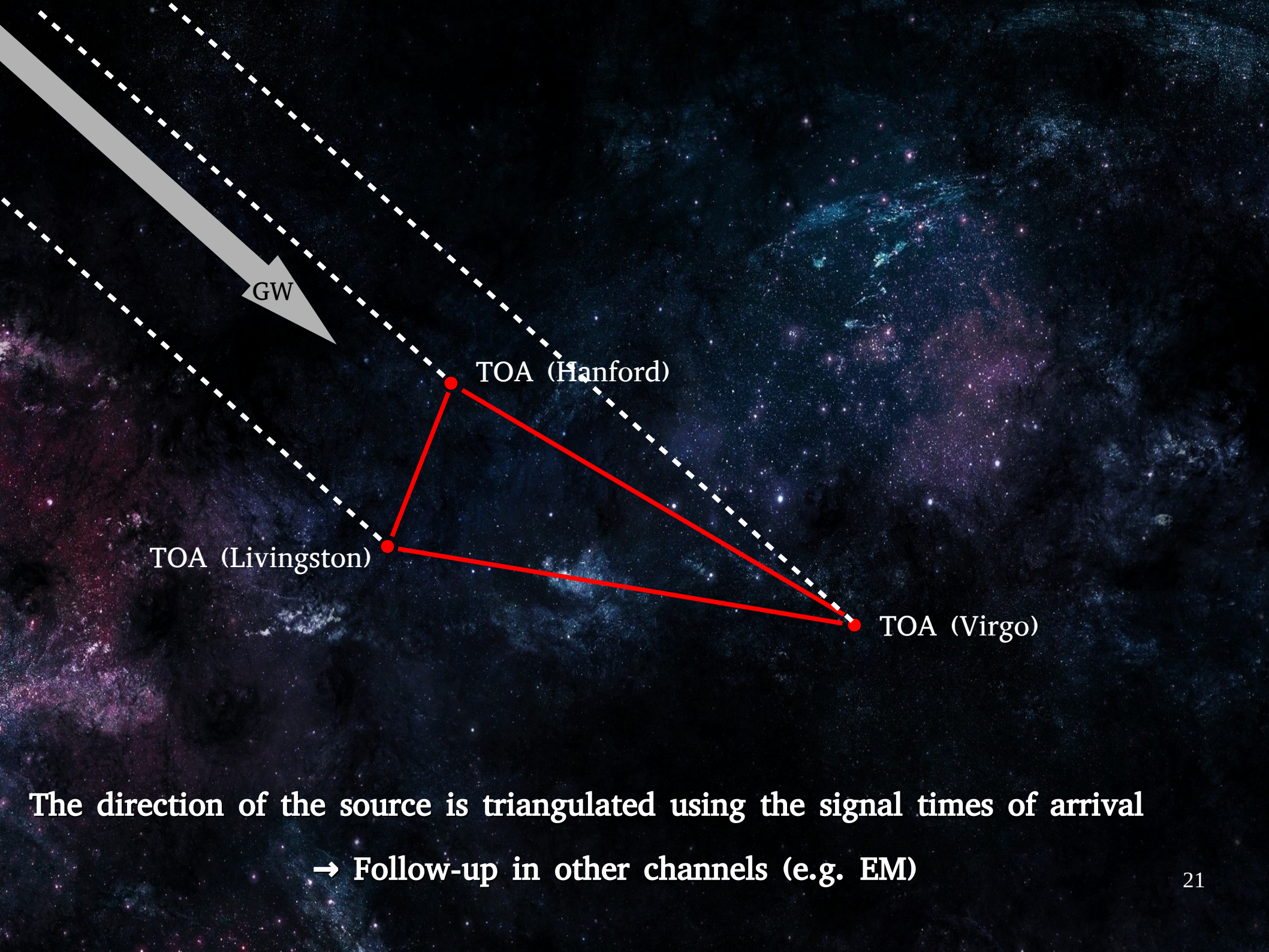
- O1 : ~50 days of data, 2 detectors
- O2 : ~100 days of data, 2(+1) detectors
- O3 : ~200 days of data, 3 detectors

Working with a network of detectors is mandatory

- to perform a coincident search
- to test the signal consistency across the network
- to estimate your background noise
- to locate the source of gravitational waves







GW

TOA (Hanford)

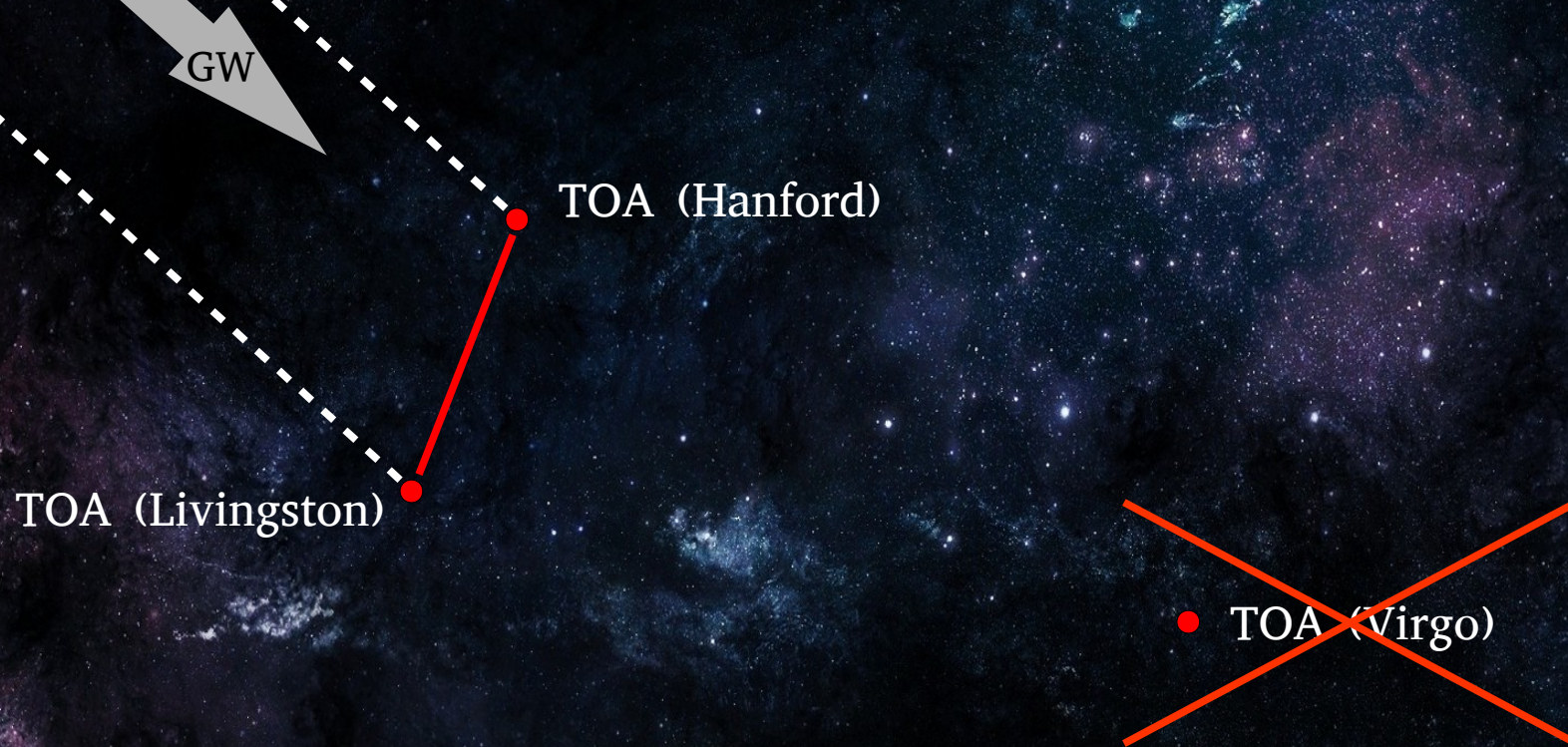
TOA (Livingston)

TOA (Virgo)

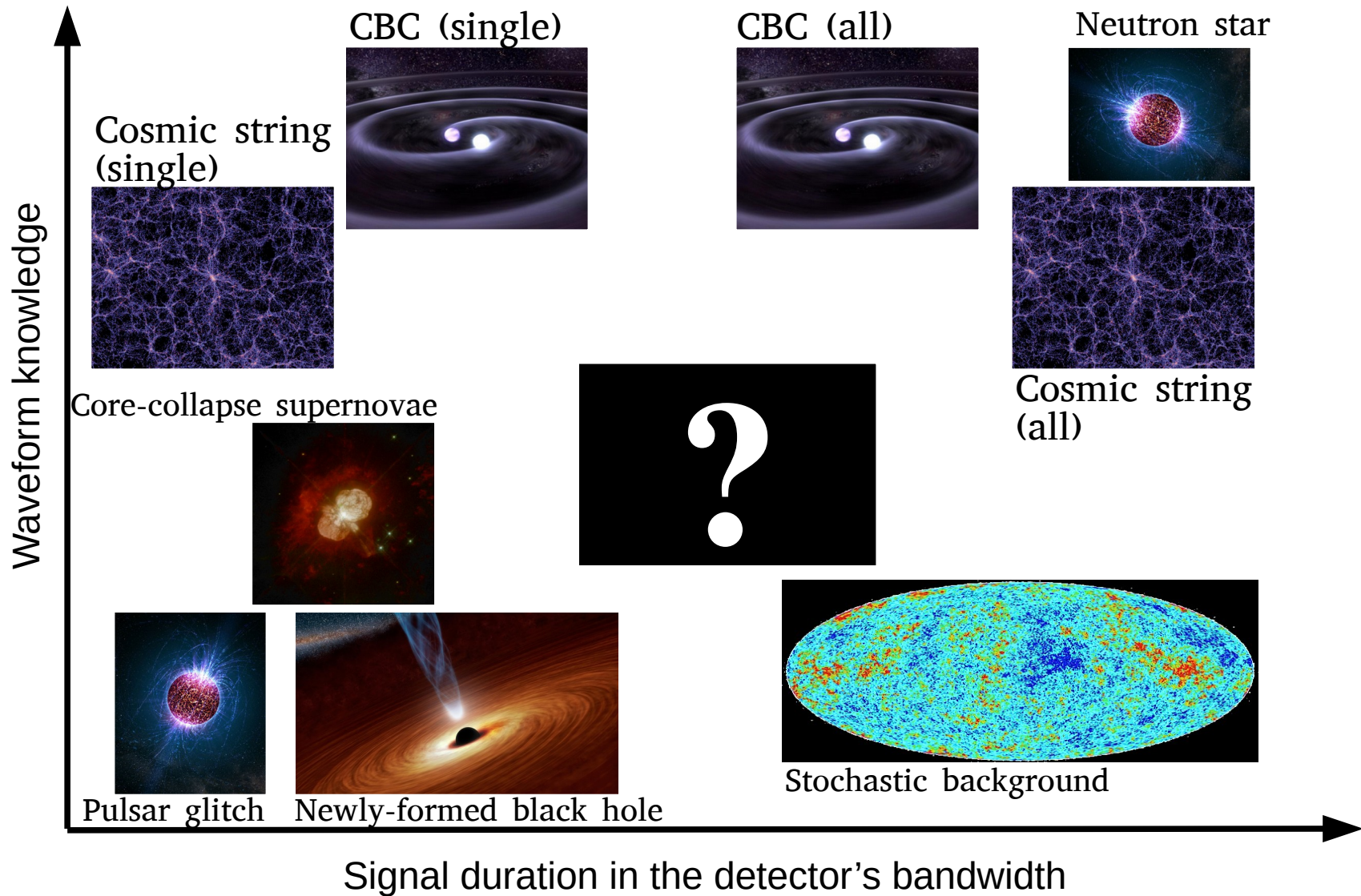
The direction of the source is triangulated using the signal times of arrival

→ Follow-up in other channels (e.g. EM)

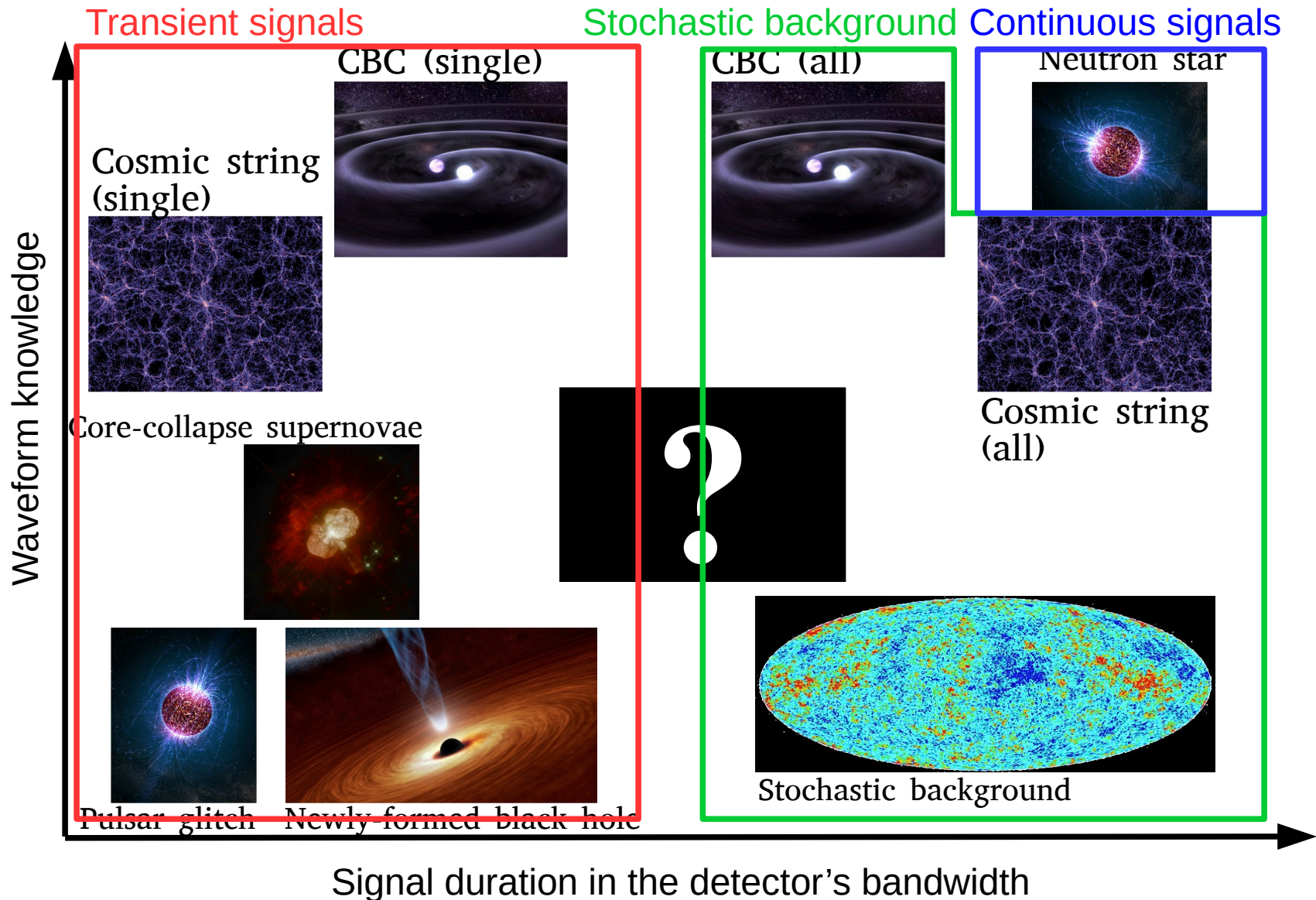
What localization to expect if the signal is detected by only 2 detectors?



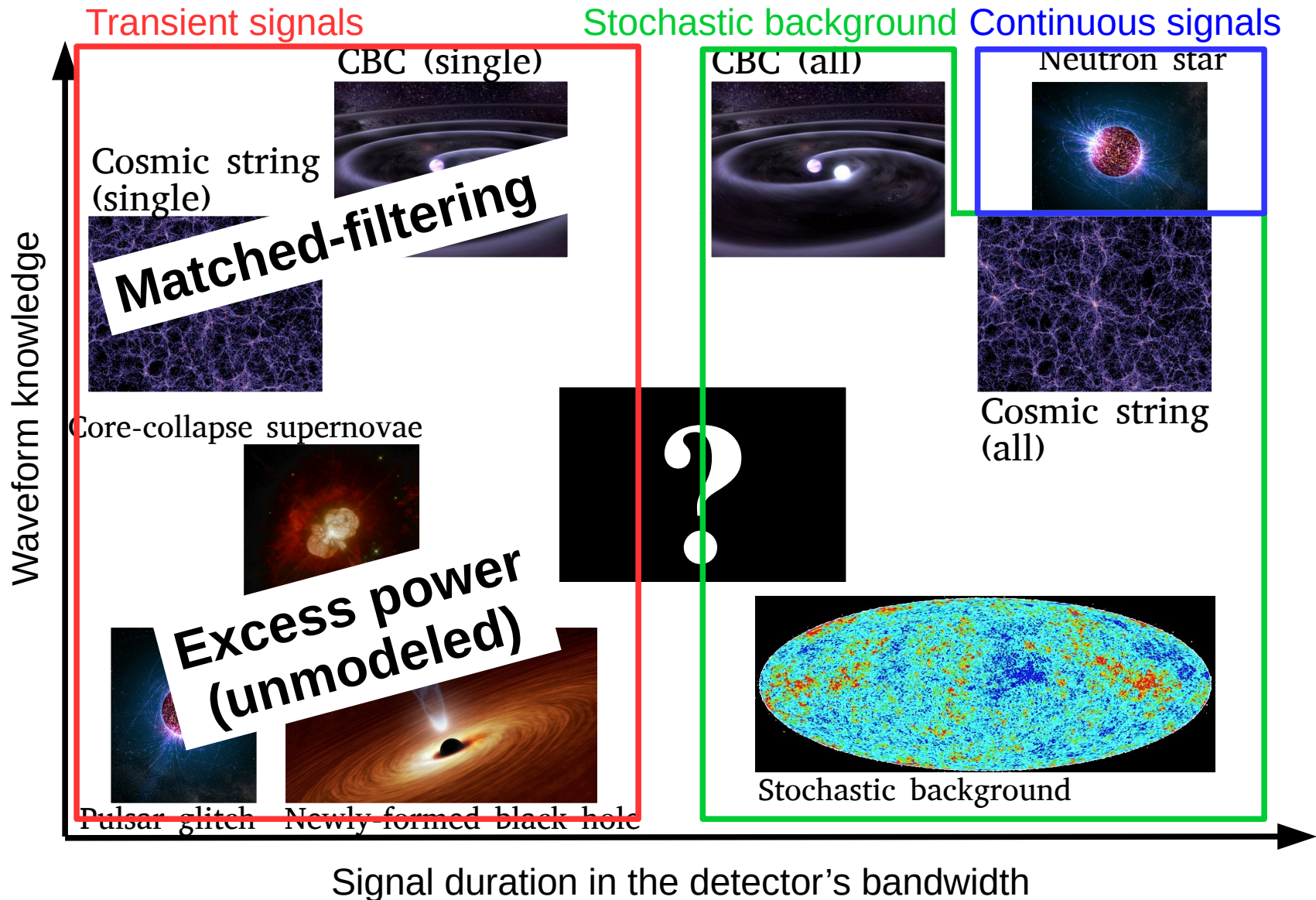
Source classification: search method



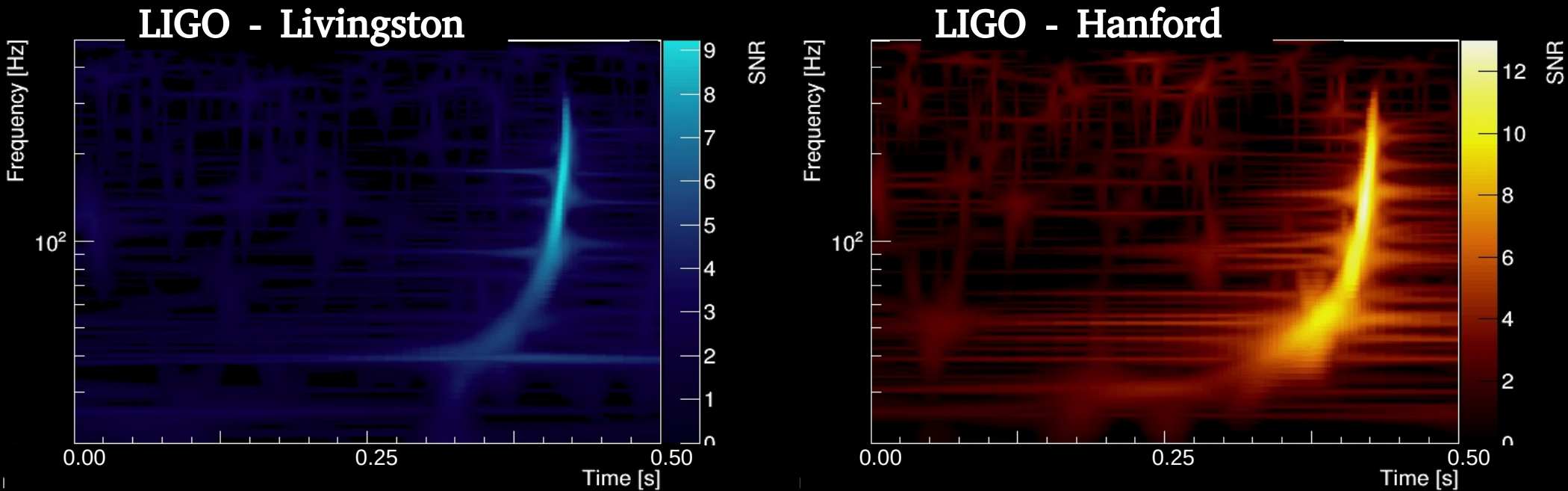
Source classification: search method



Source classification: search method



Unmodeled searches



- Time-frequency decomposition → noise events + (maybe) GW events
- Coincidence between detectors (time + other parameters)
- Noise rejection
- Classify events using a “smart” recipe
- Estimate background
- Compare events with your background
- Measure the probability for each events to be true GW signal

GW150914

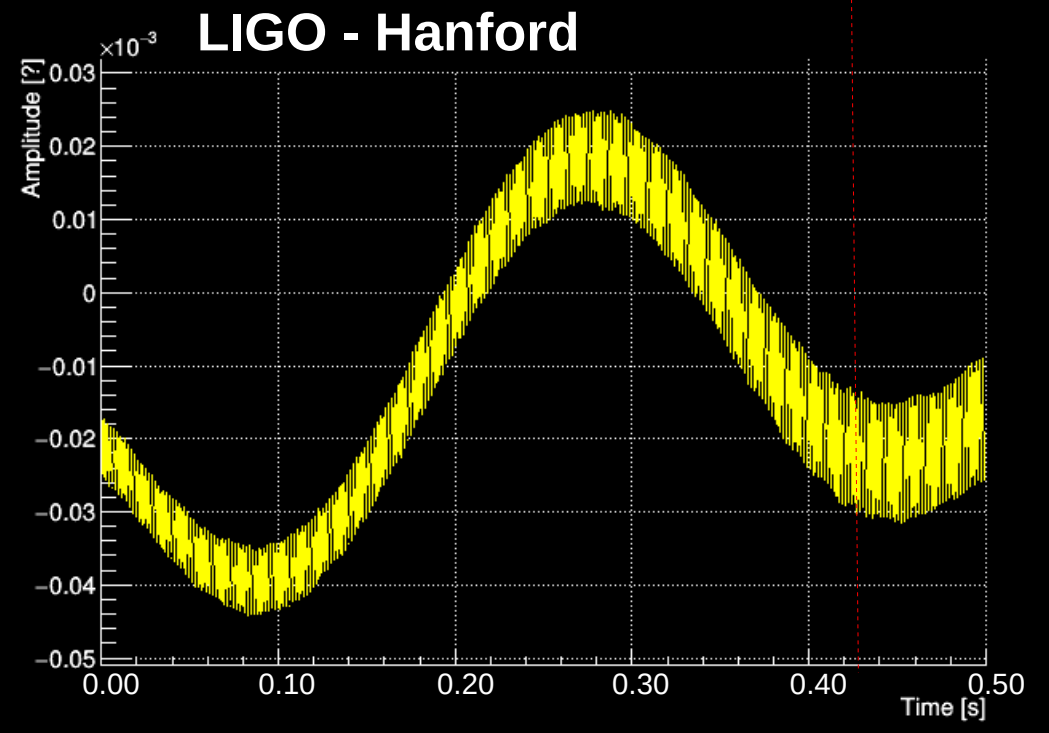
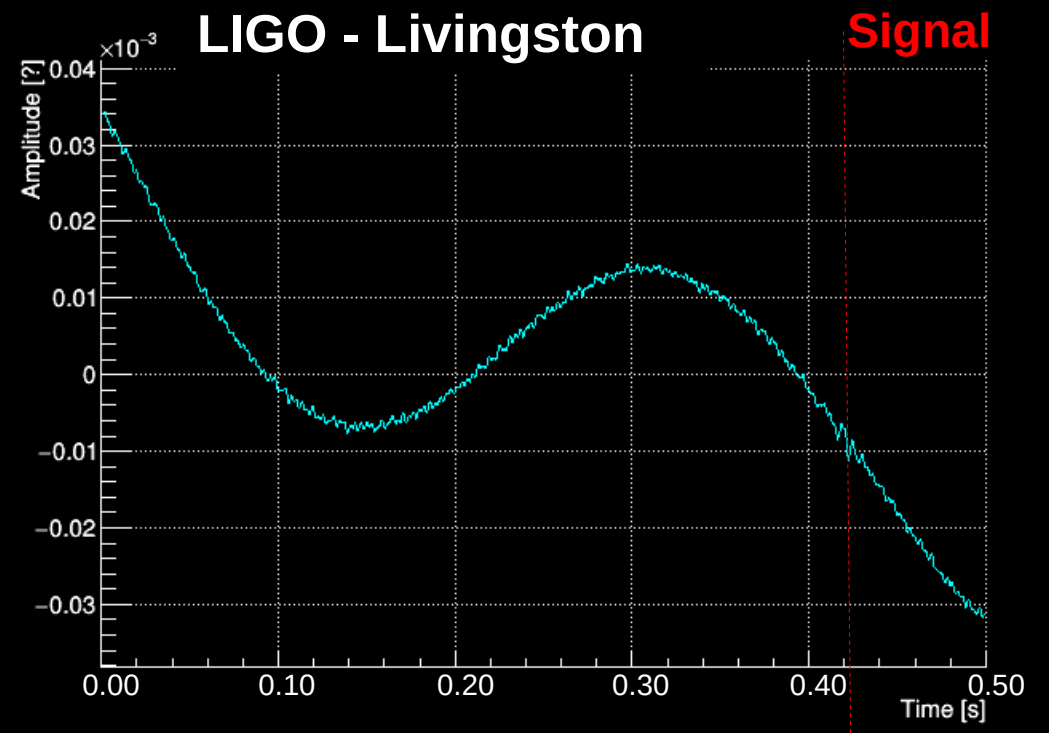
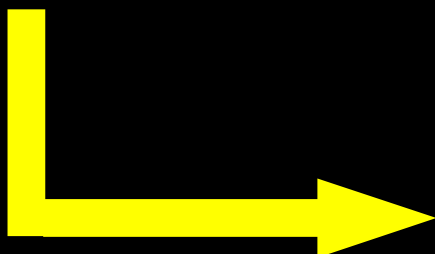
Output power



Livingston

Hanford

Output power

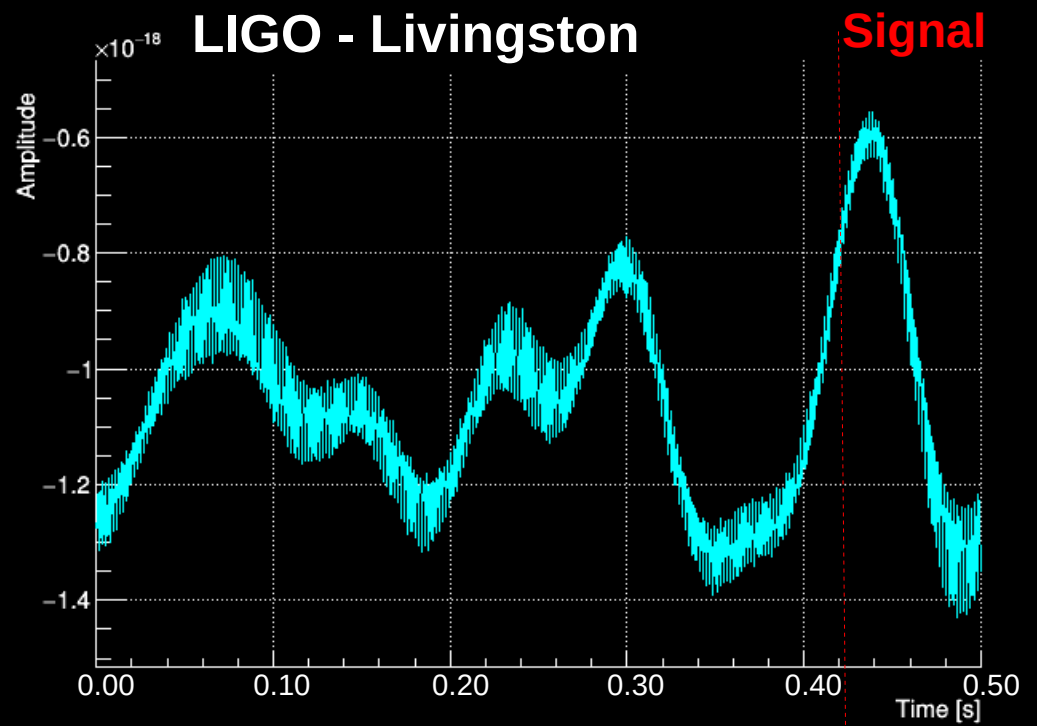


GW150914

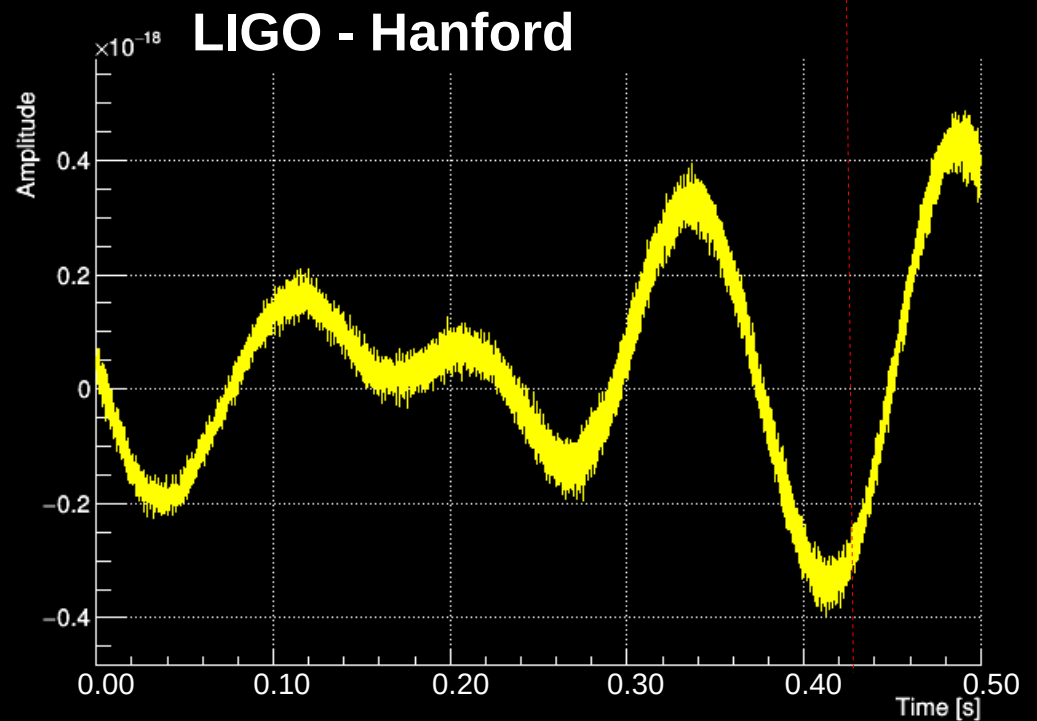
$$h(t)$$

Data is calibrated

→ GW strain amplitude $h(t)$
(including high-pass filter $f > 10$ Hz)

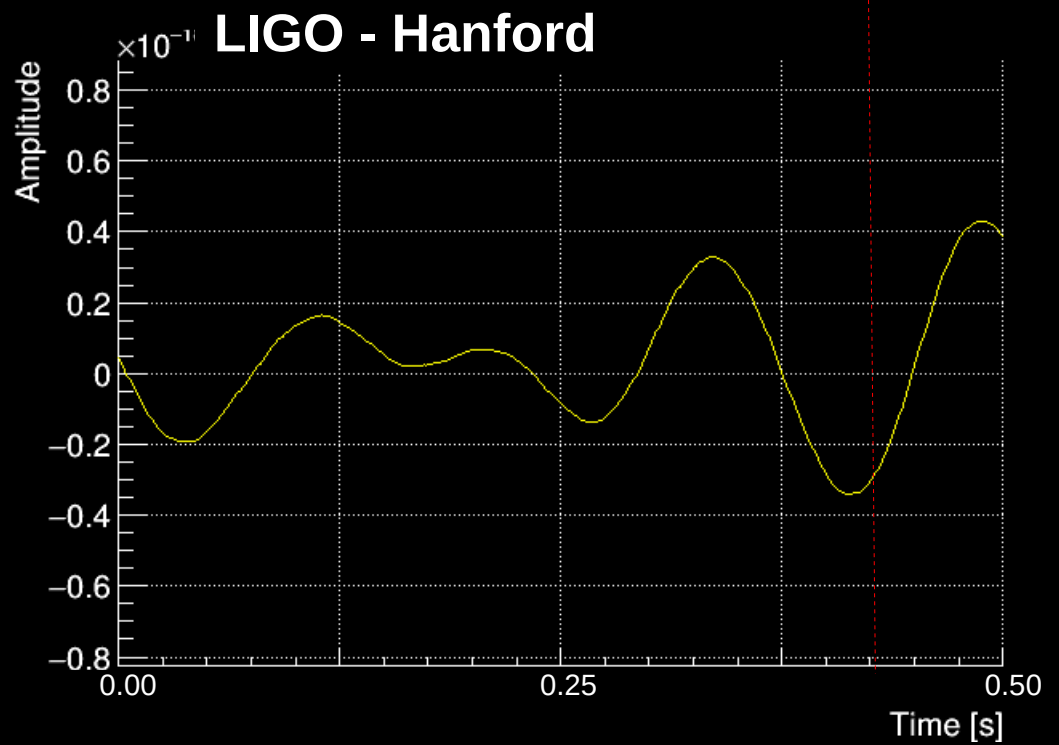
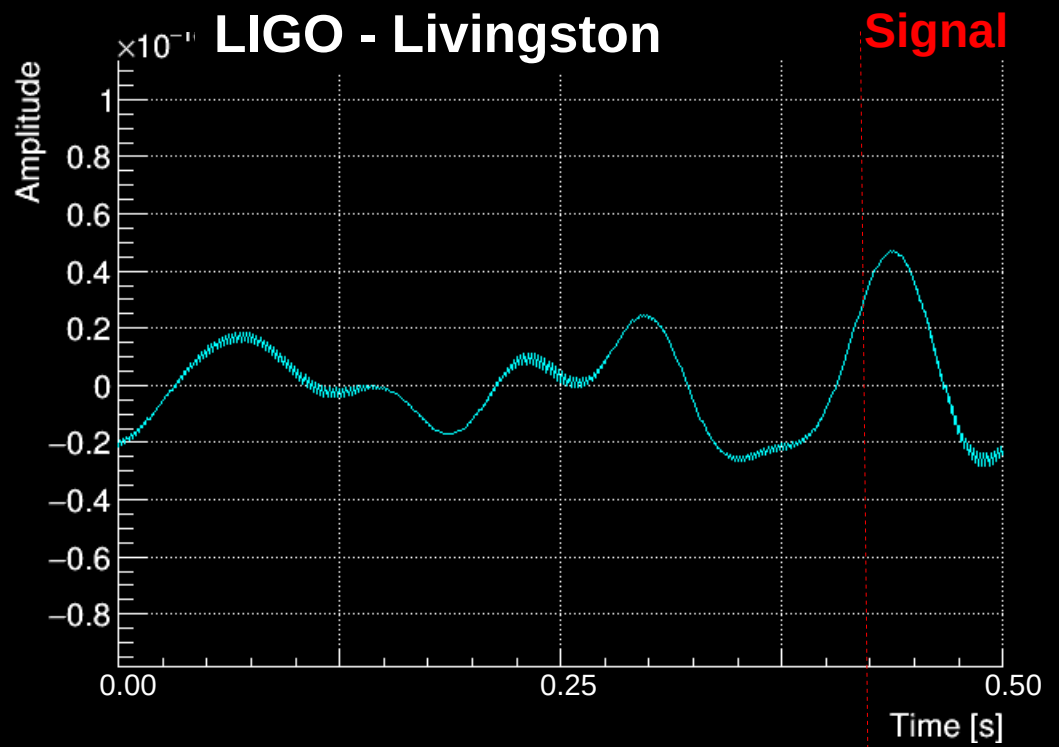


$$h(t)$$



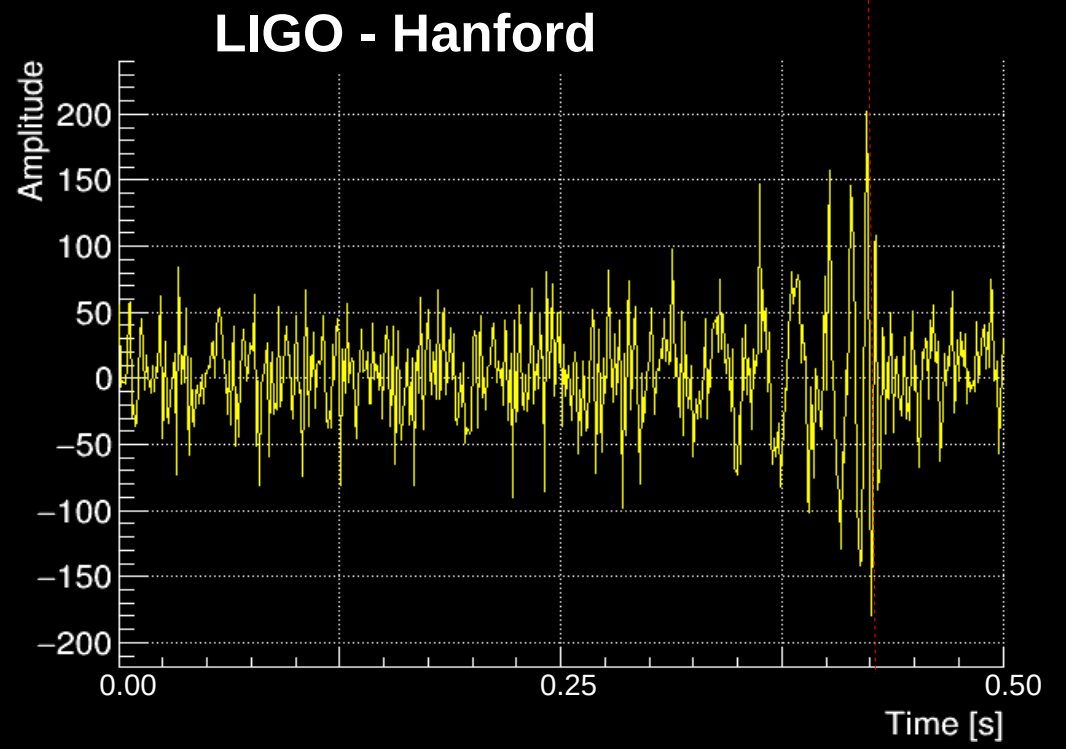
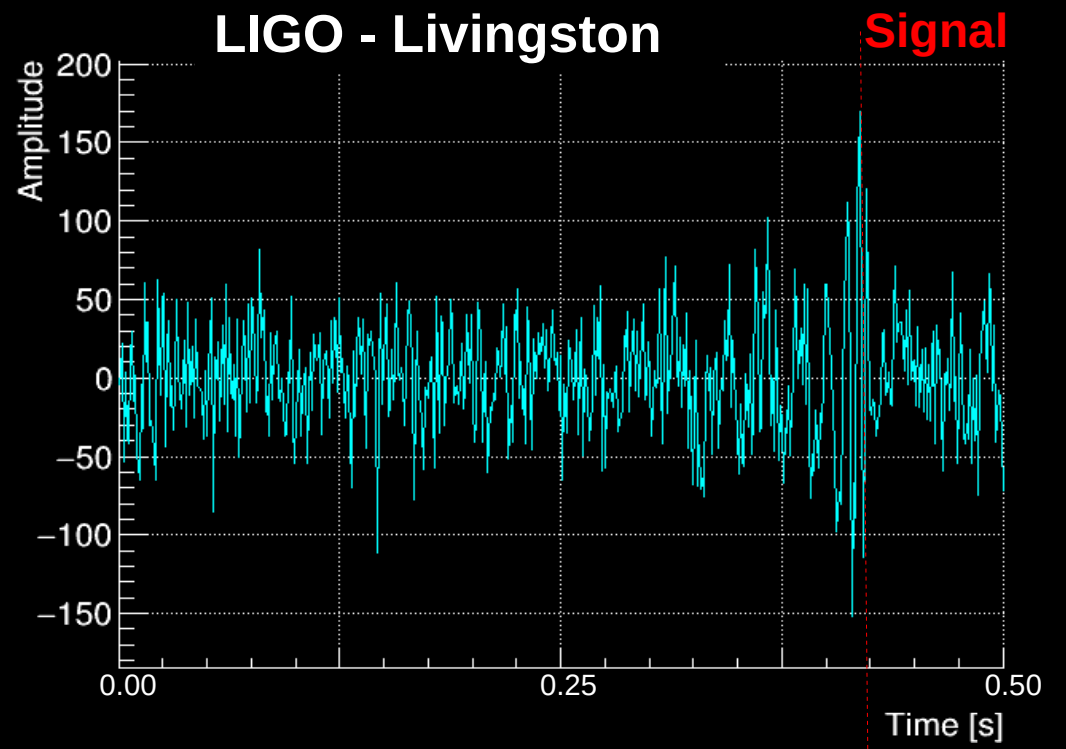
GW150914

Data are low-pass filtered
(here, < 500 Hz)



GW150914

Data are whitened

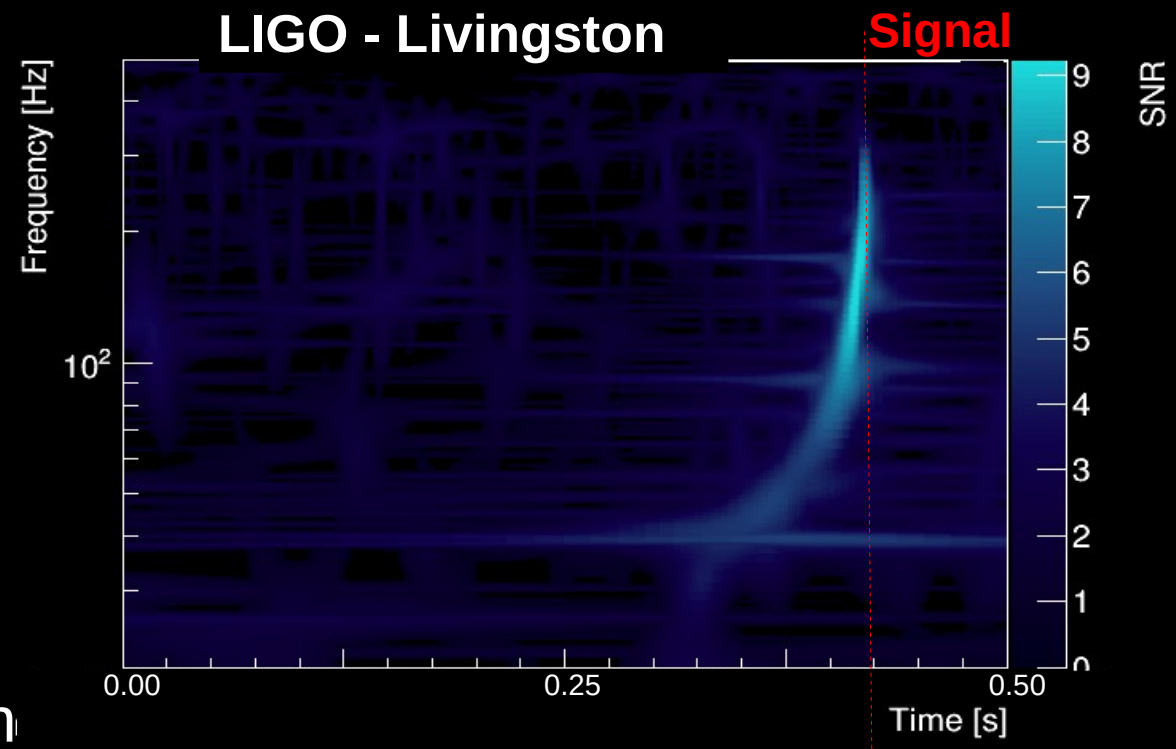


GW150914

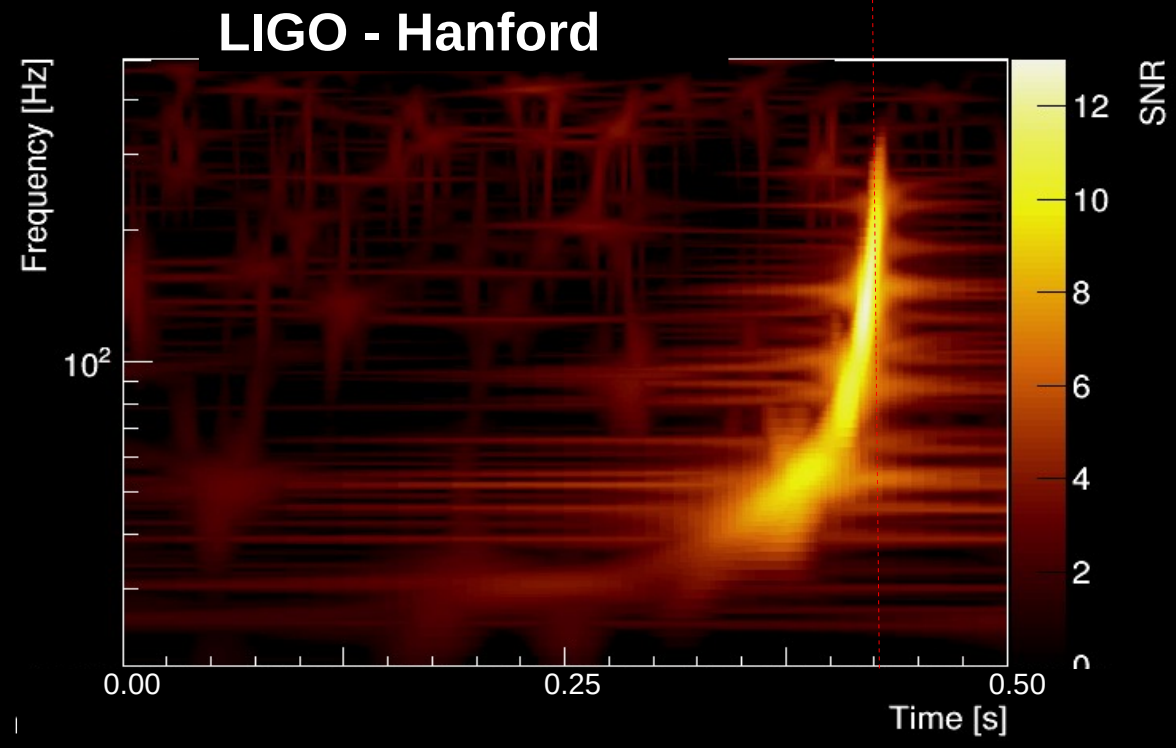
Time-frequency decomposition
(Short Fourier transforms)

$$X(\tau, \phi, Q) = \int_{-\infty}^{+\infty} h_{det}(t) w(t - \tau, \phi, Q) e^{-2i\pi\phi\tau} dt$$

LIGO - Livingston



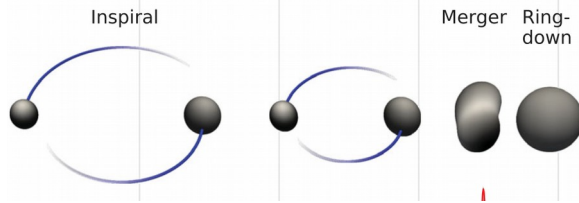
LIGO - Hanford



Modeled search

Theoretical input:

- 90s: CBC PN waveforms (Blanchet, Iyer, Damour, Deruelle, Will, Wiseman, ...)
- 00s: CBC Effective One Body “EOB” (Damour, Buonanno)
- 06: BBH numerical simulation (Pretorius, Baker, Loustos, Campanelli)



The intrinsic waveform parameters:

- Masses:

$$M_{tot} = M_1 + M_2$$

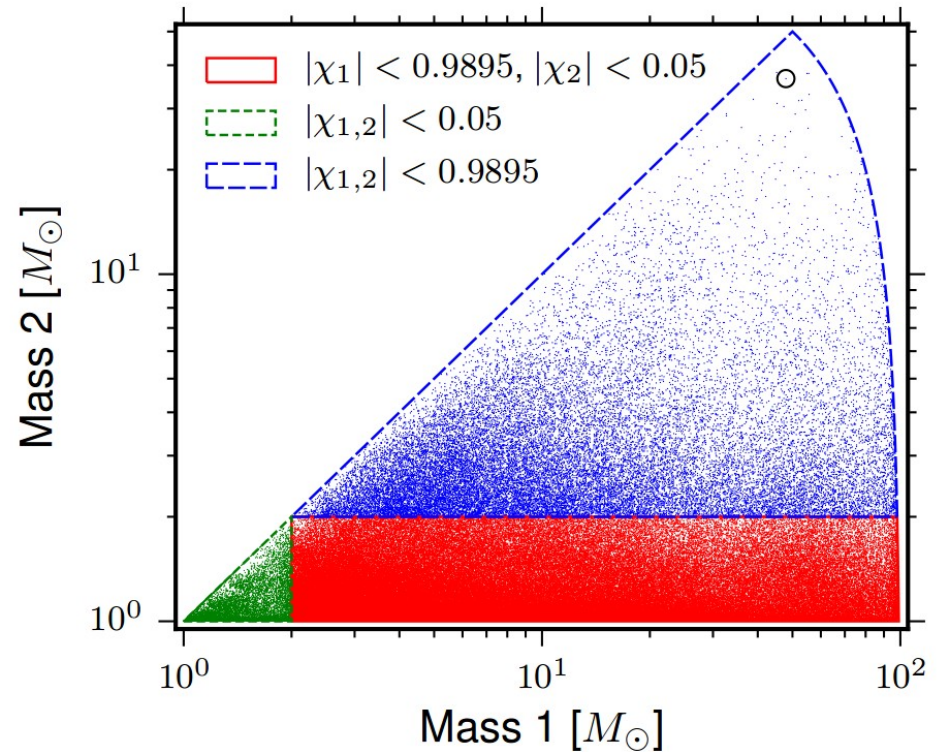
- Spins and orbital angular momentum:

$$\vec{S}_{tot} = \vec{S}_1 + \vec{S}_2 + \vec{J}$$

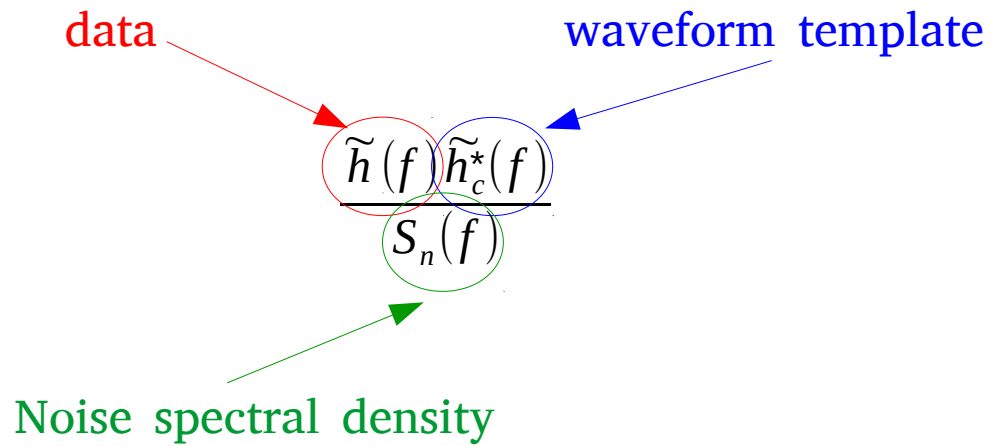
The waveform models used for the search:

- Inspiral, PN3.5 for $M_{tot} < 4 M_{sun}$
- Inspiral/Merger/Ringdown EOB + numerical relativity for $M_{tot} > 4 M_{sun}$
- Spins and orbital angular momentum are aligned

Template bank → match-filtering technique



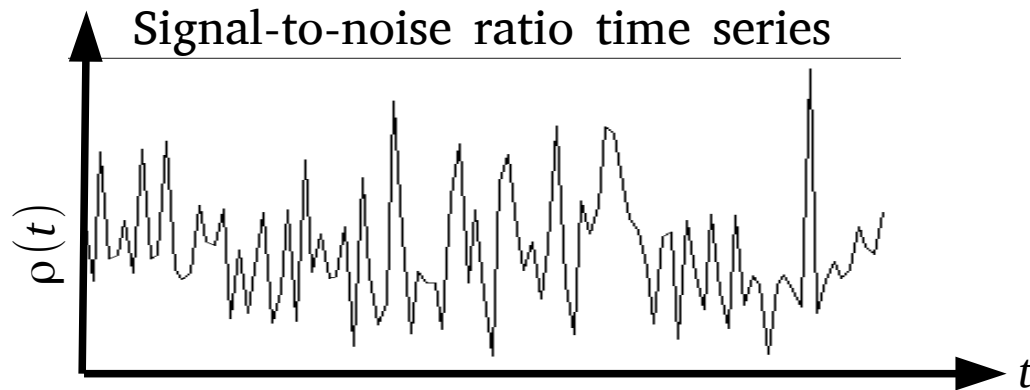
Match filtering



Match filtering

$$\rho_c(t) = 4 \Re \left[\int_0^\infty \frac{\tilde{h}(f) \tilde{h}_c^*(f)}{S_n(f)} e^{2i\pi f t} df \right]$$

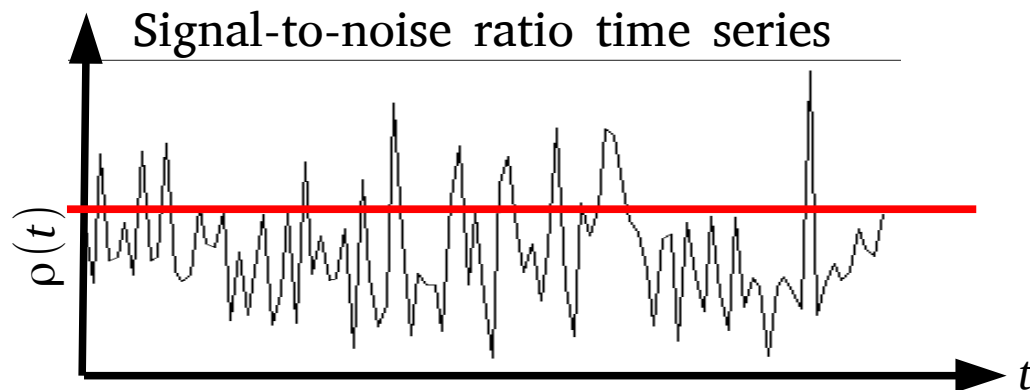
A signal-to-noise ratio (SNR) is computed for each template



Match filtering

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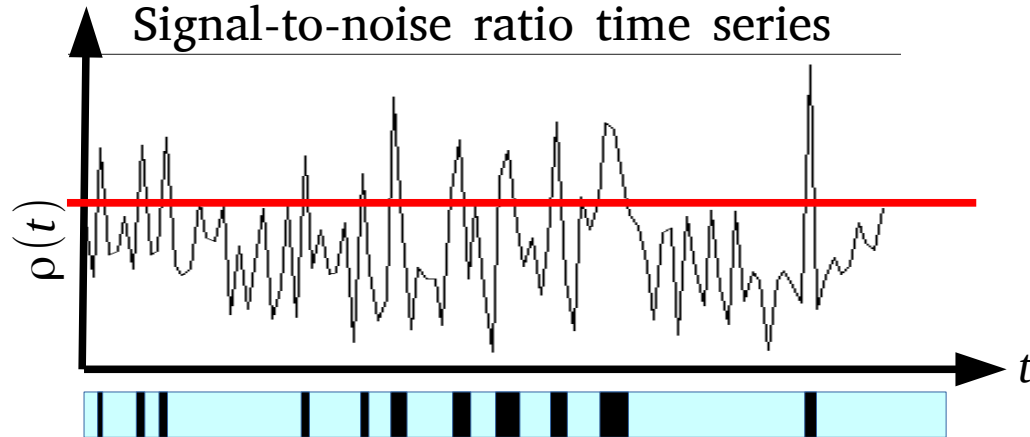


A list of events is produced:

- start/end/peak times
- SNR
- template parameters (masses, spins)

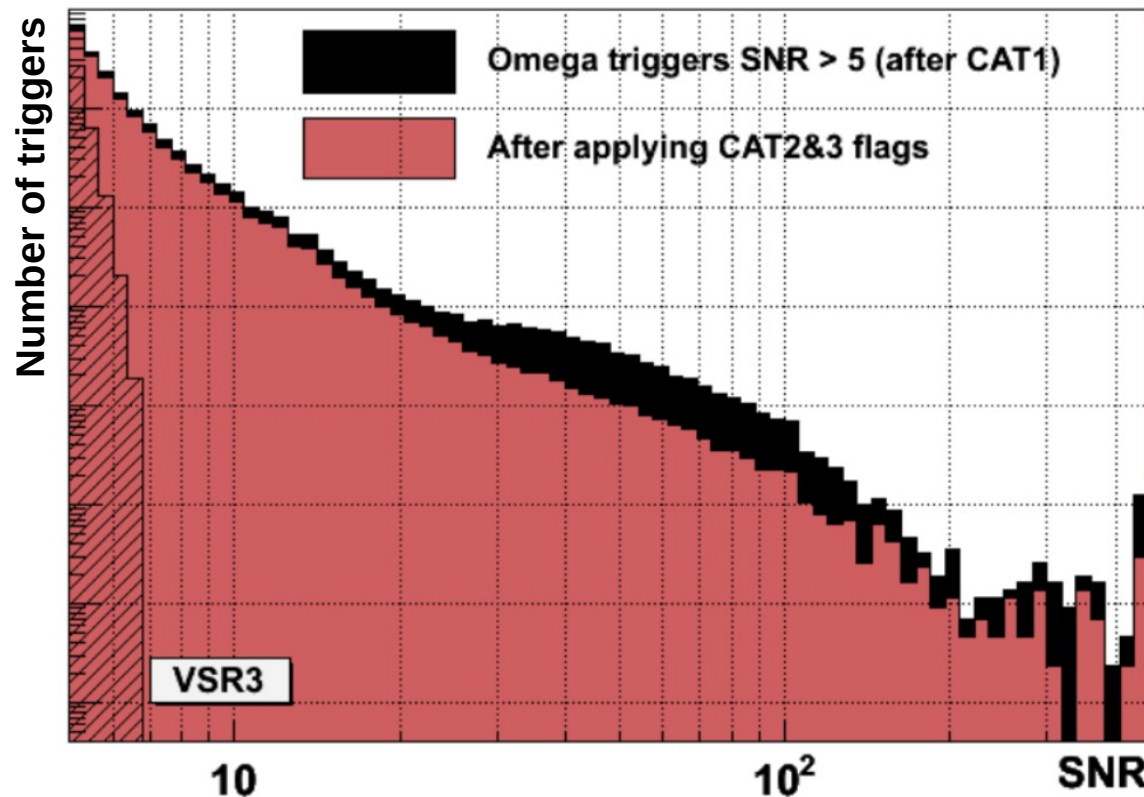
Now the challenge is to reject noise events to better isolate true signals

Single-detector triggers



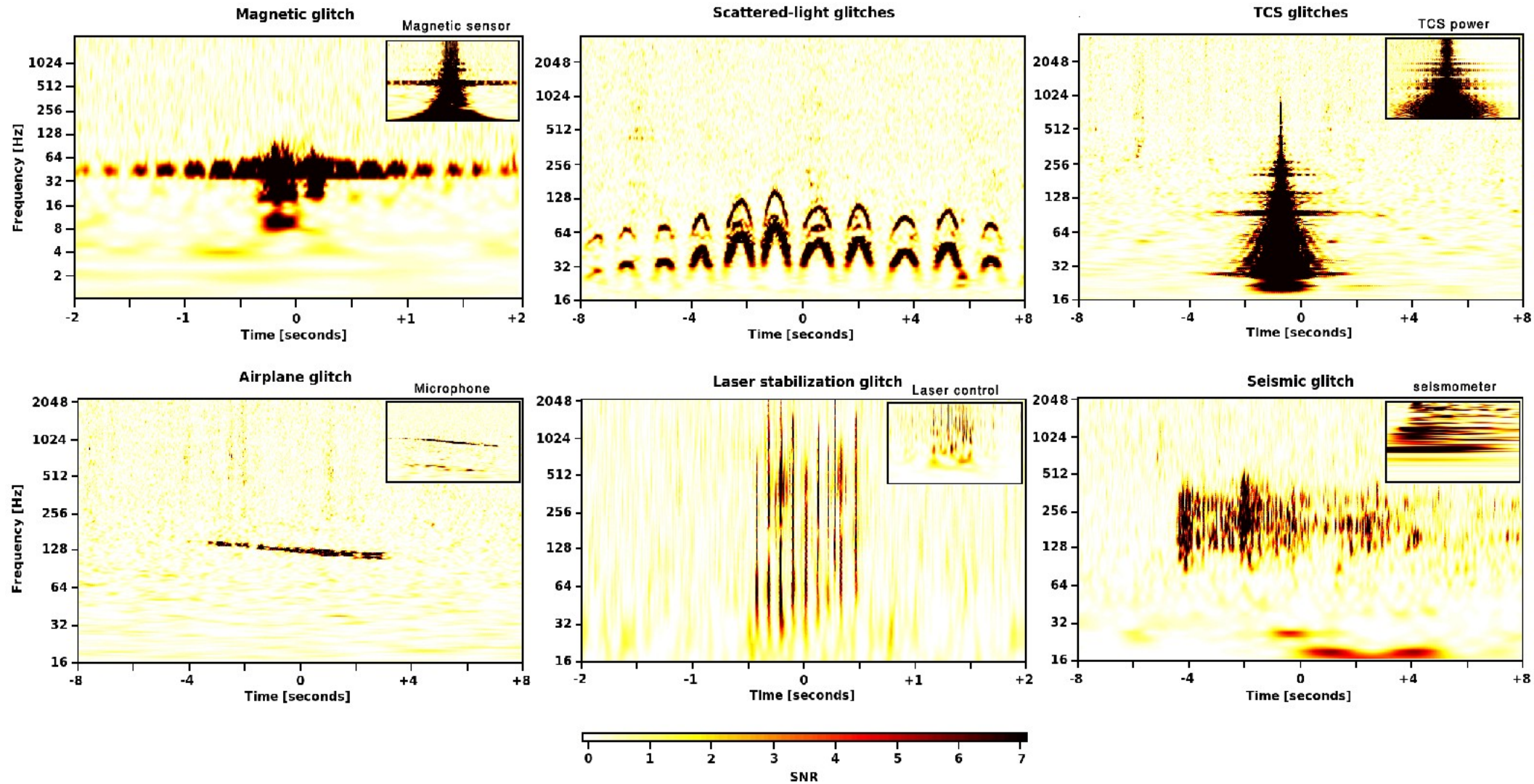
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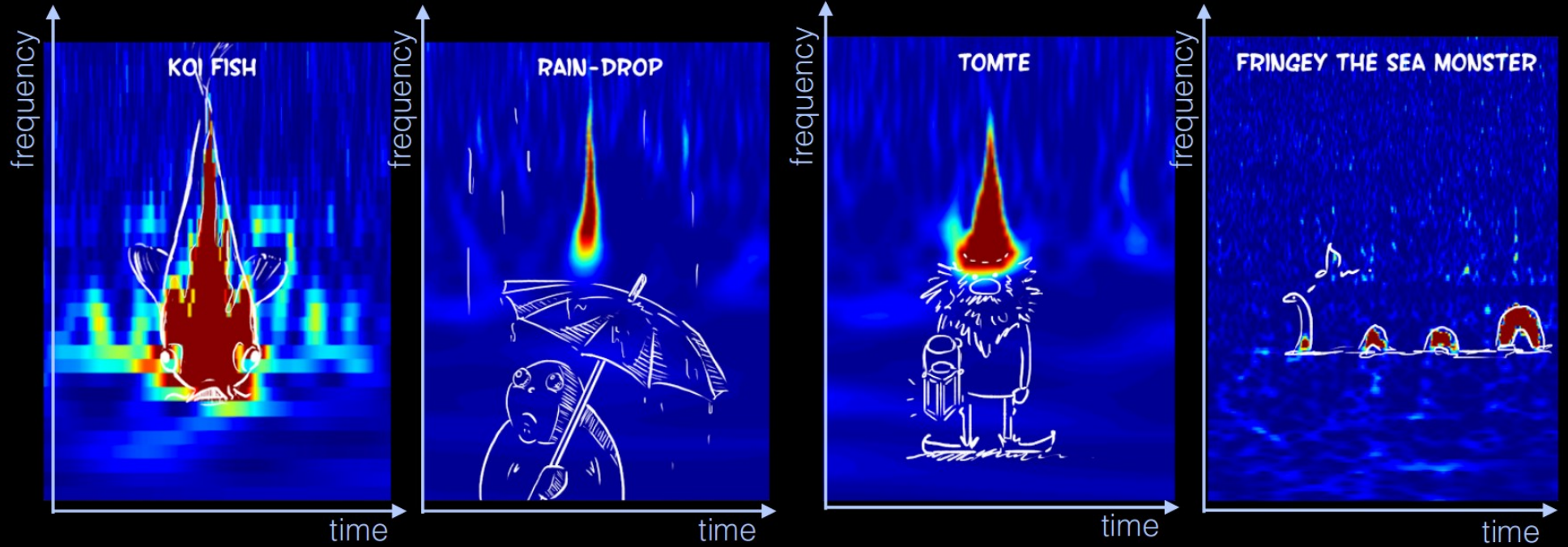
The noise distribution is highly non-Gaussian !

Single-detector noise



Single-detector noise

THE ART OF NAMING GLITCHES



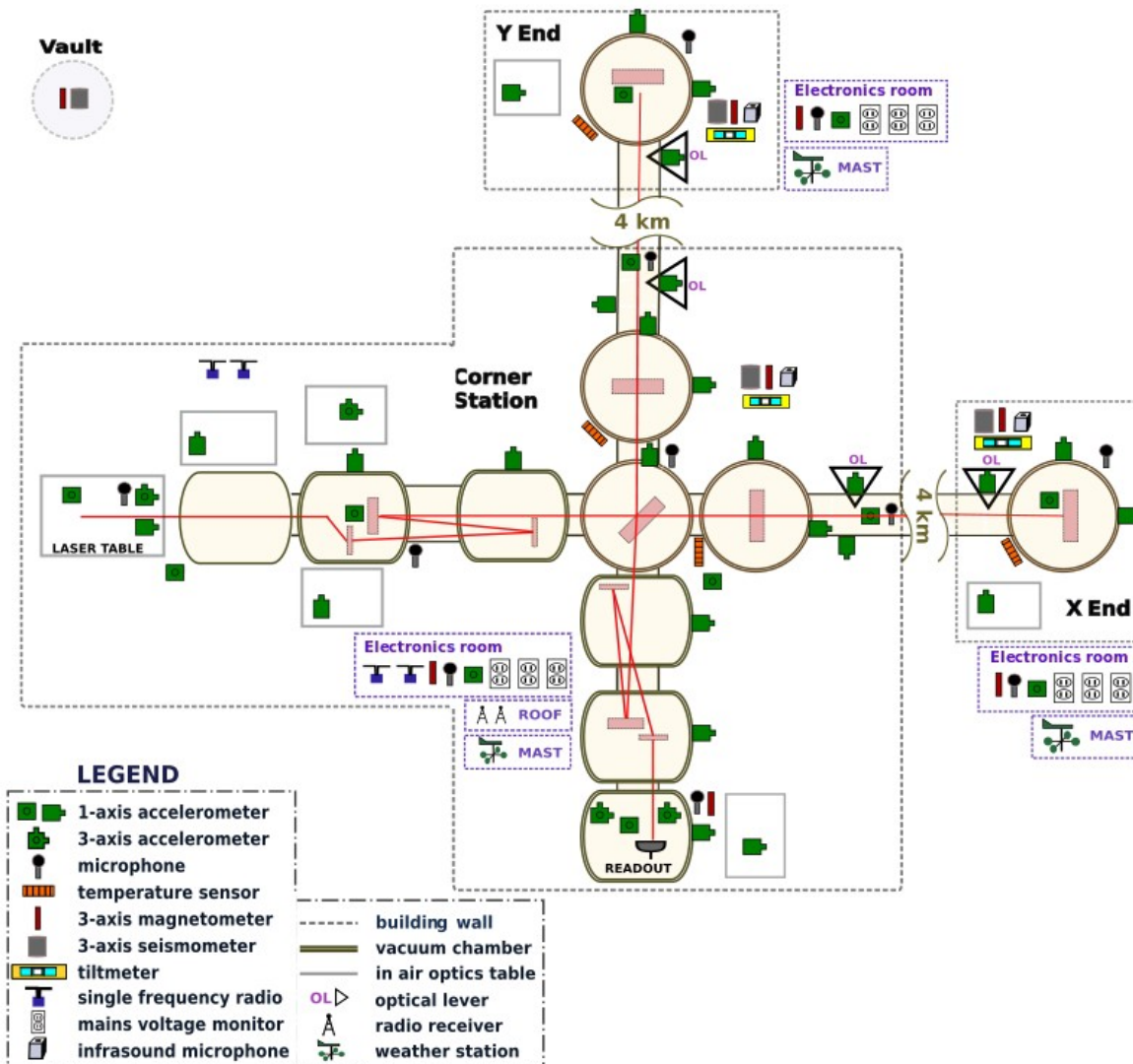
Nutsinee Kijbunchoo

ANTIMATTERWEBCOIMICS.COM

Monitoring noise

Thousands of auxiliary channels are used to monitor the instruments

- environmental sensors
- detector sub-systems
- detector control



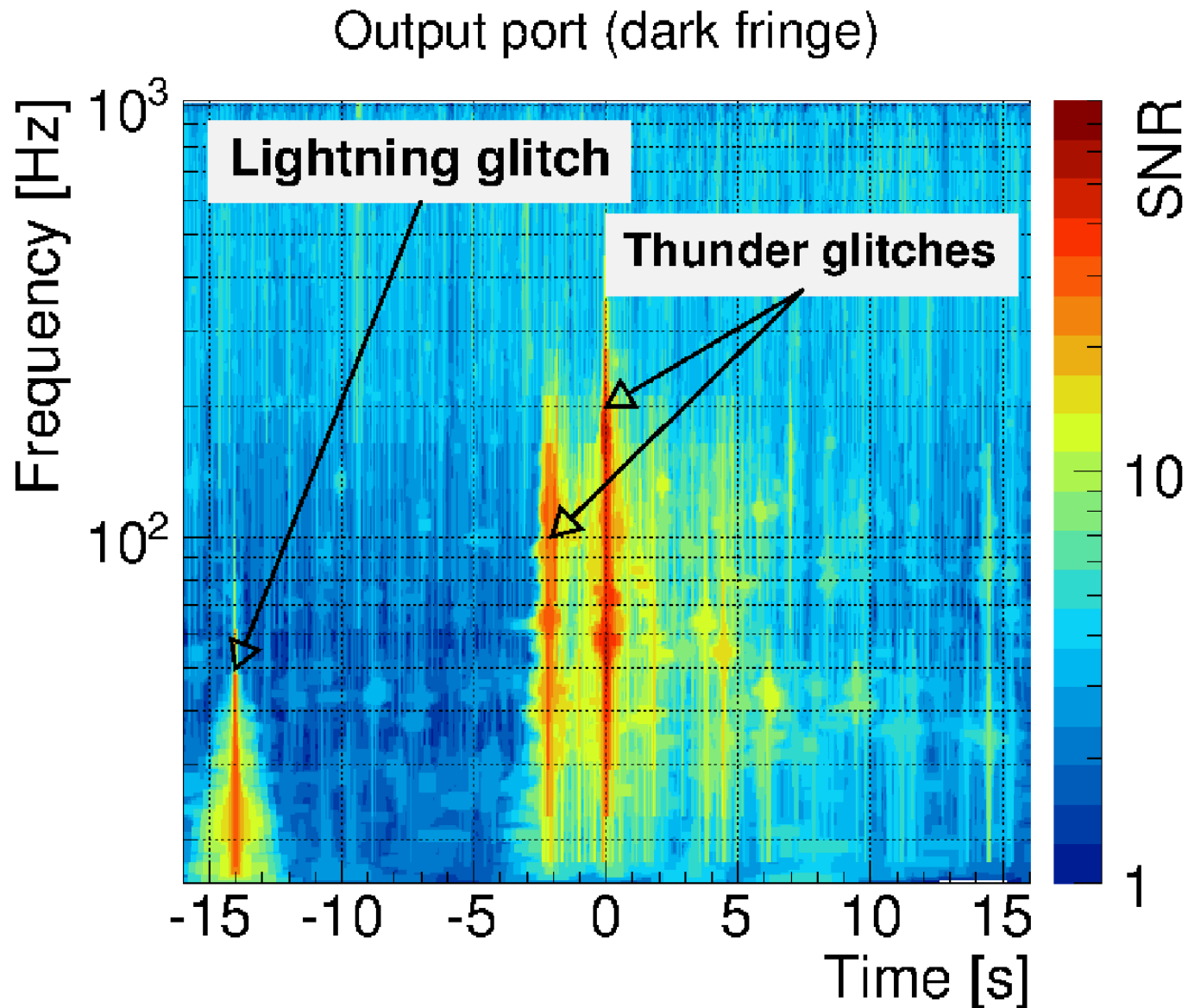
Noise injection campaigns are conducted to identify the detector's response to different noise stimulation

Multiple transient noises were identified during the run

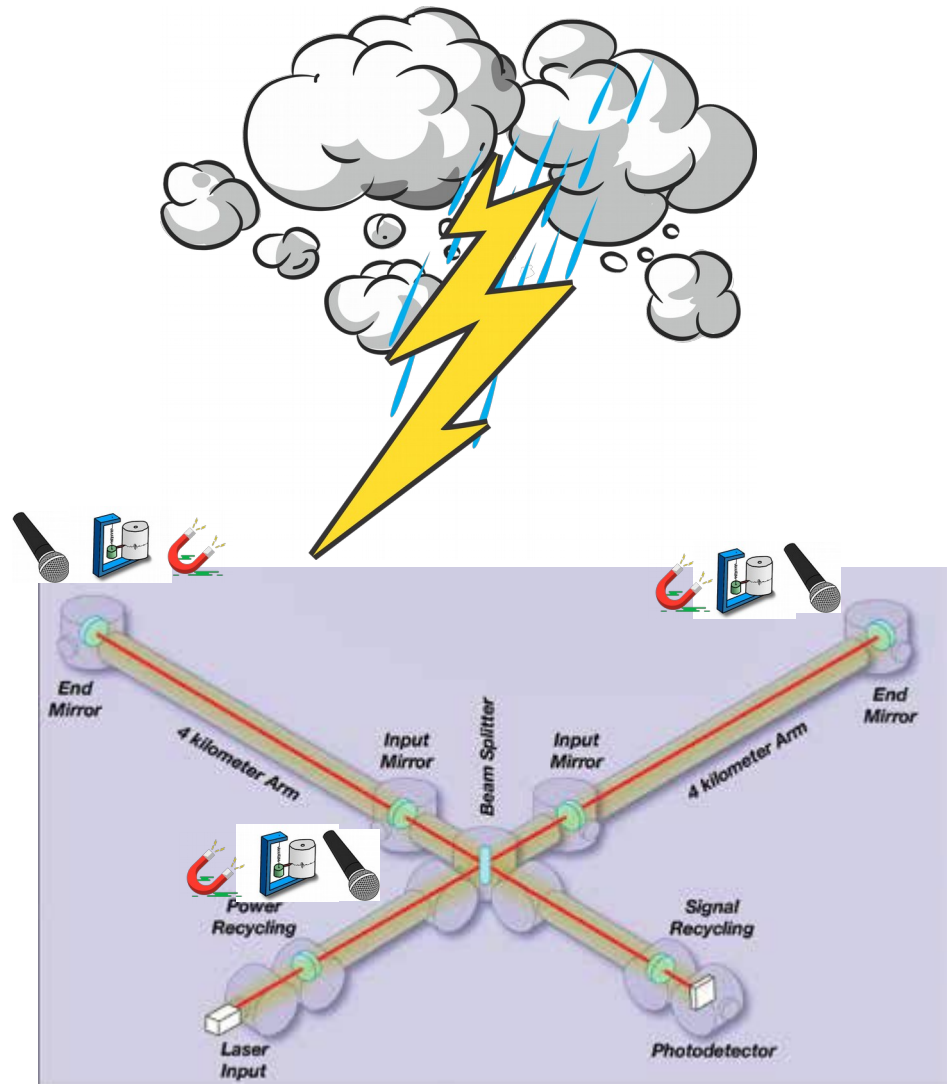
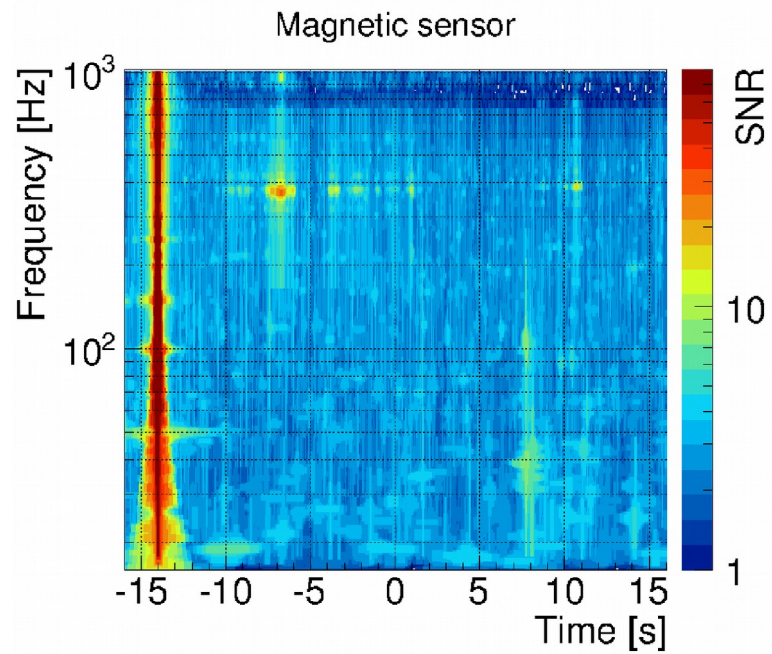
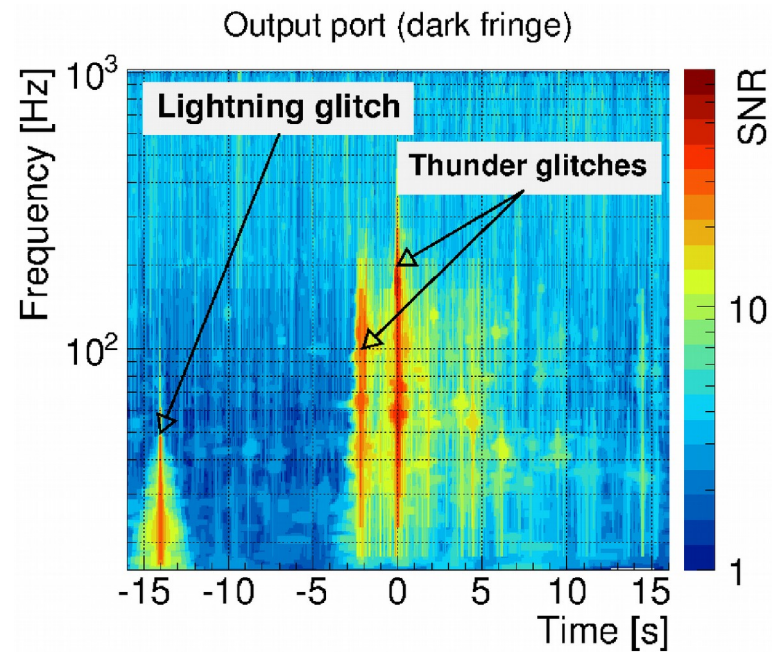
- Anthropogenic noise
- Earthquakes
- Radio-frequency modulation
- ...

Option #1: fix the detector
 Option #2: remove transient events in the data

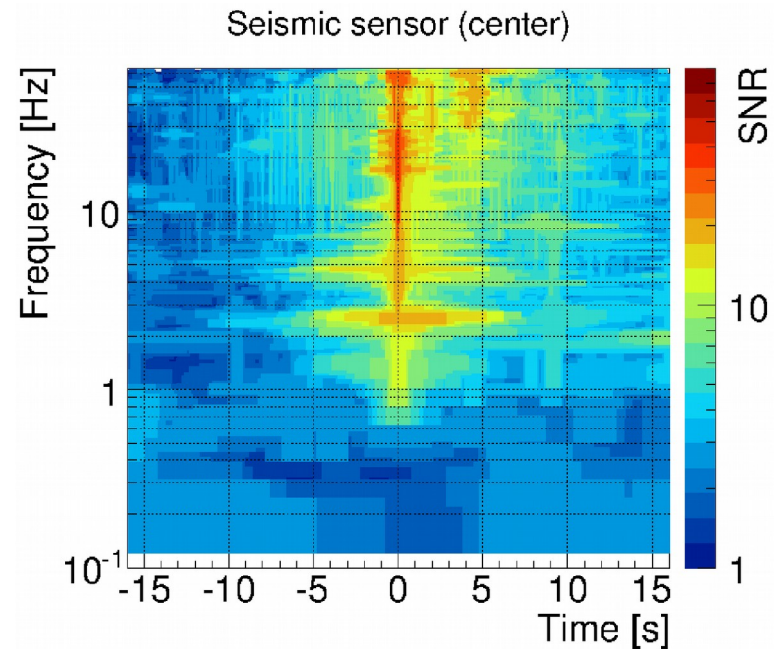
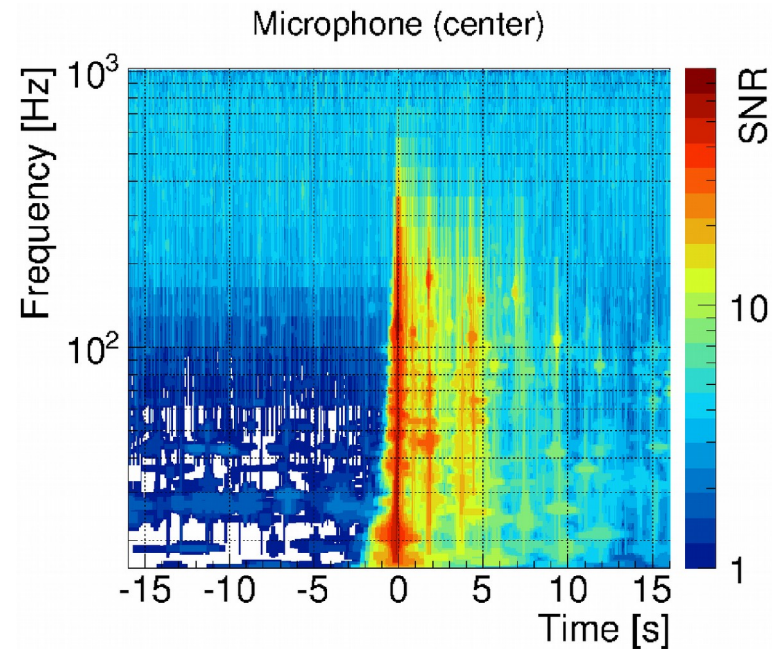
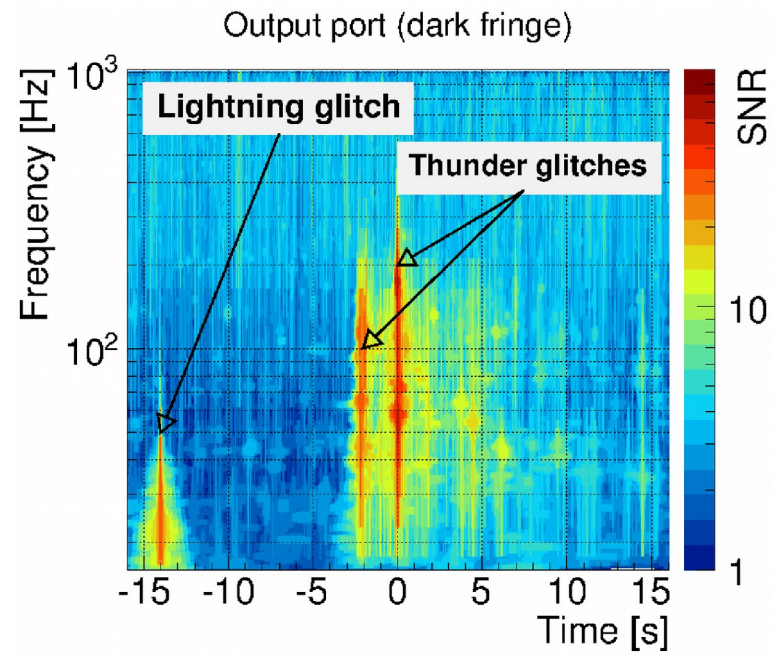
Monitoring noise



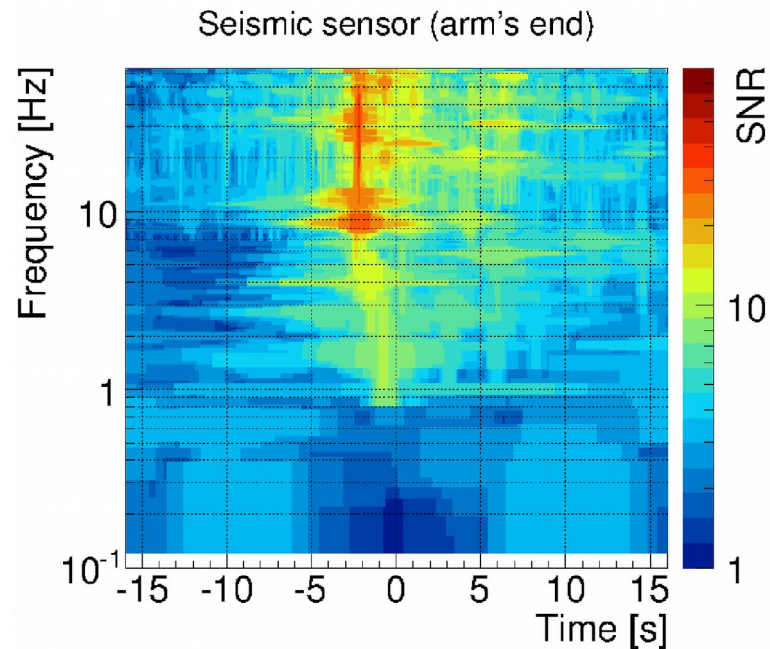
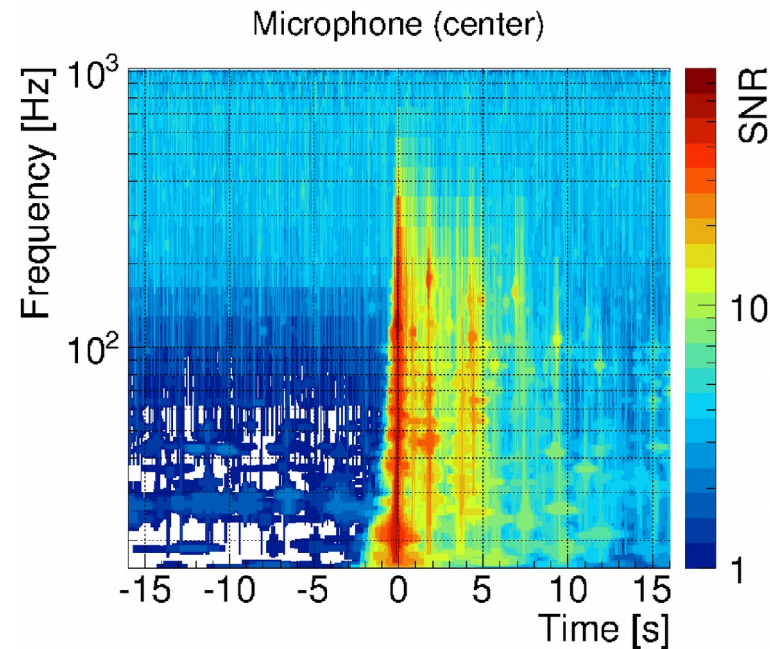
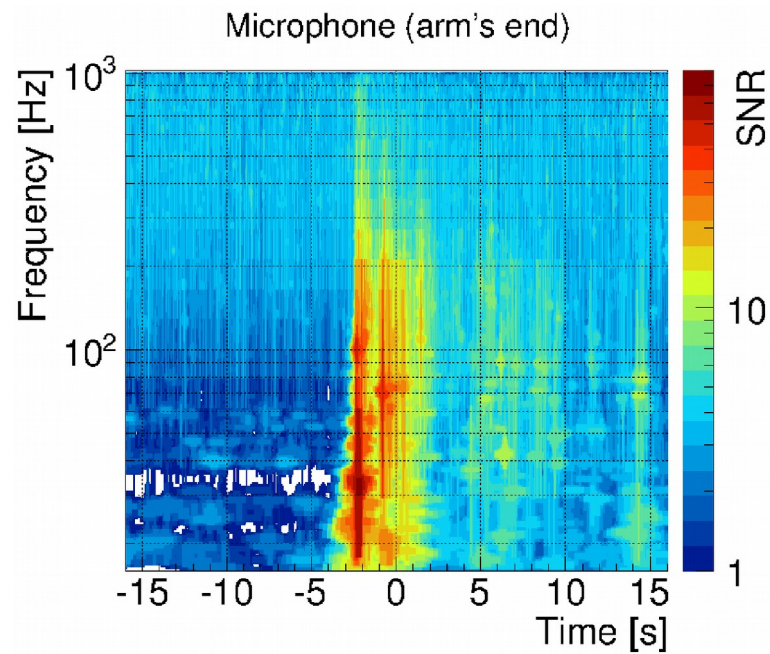
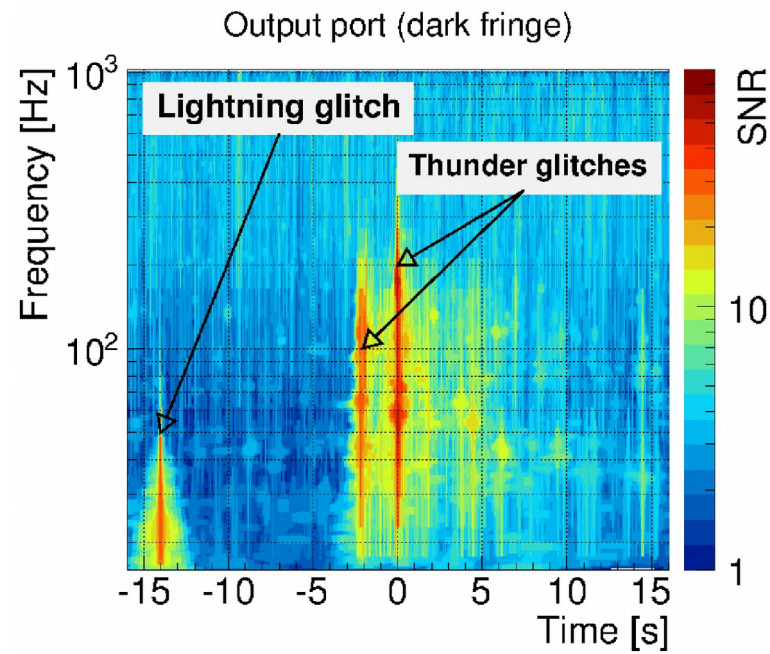
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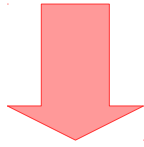
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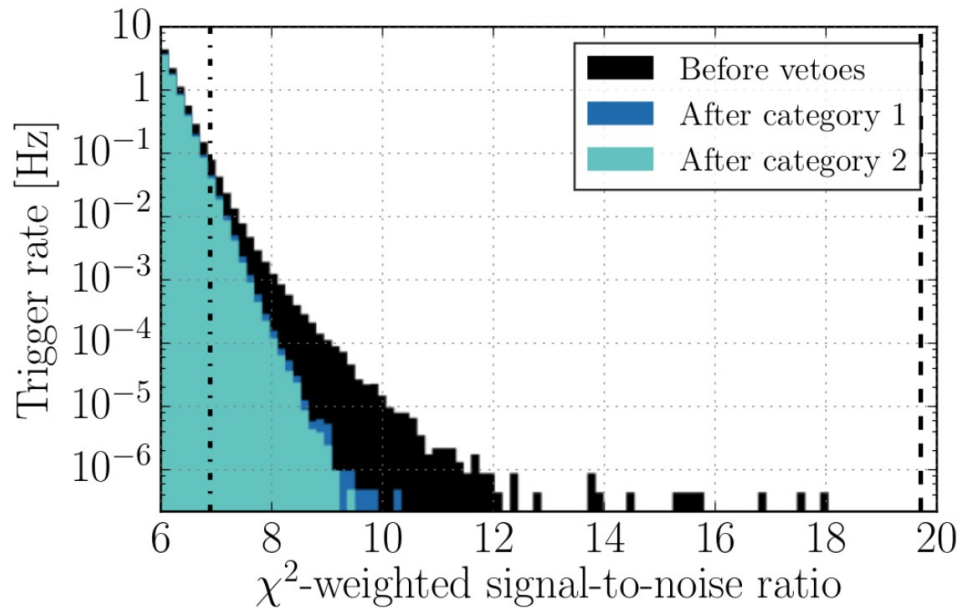
Monitoring noise



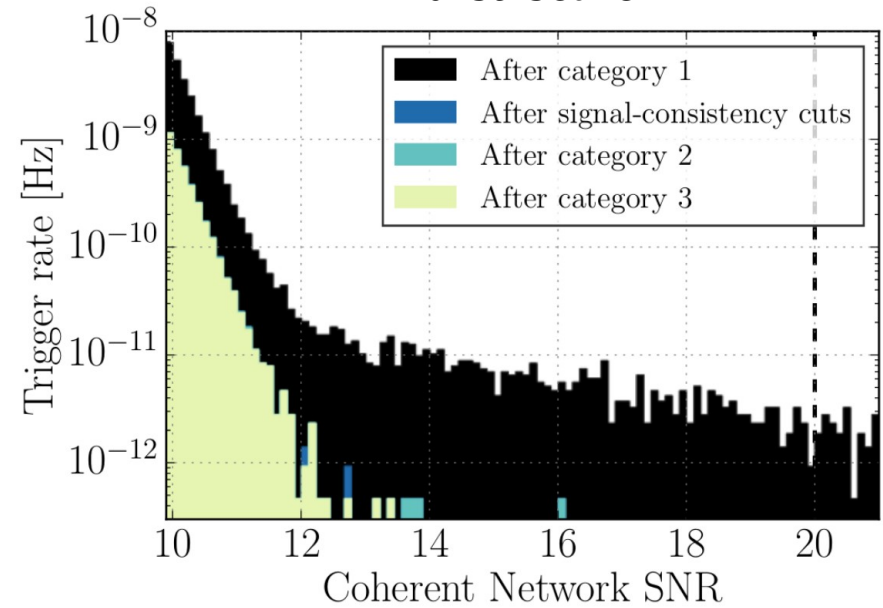
Rejecting noise



CBC search



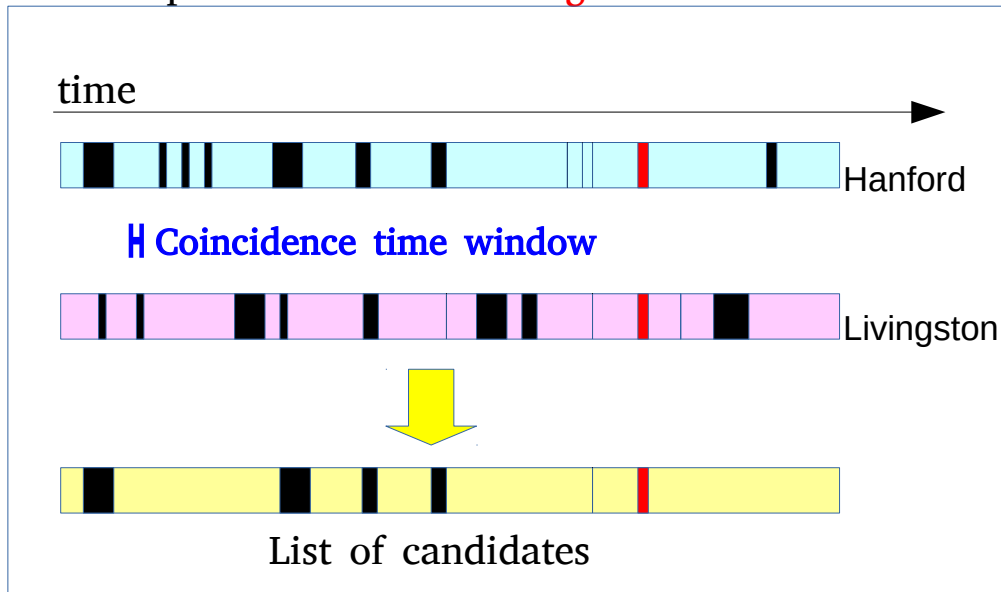
Burst search



Multi-detectors coincidence

A gravitational-wave signal is detected by multiple detectors almost simultaneously

True experiment = noise + **signal**



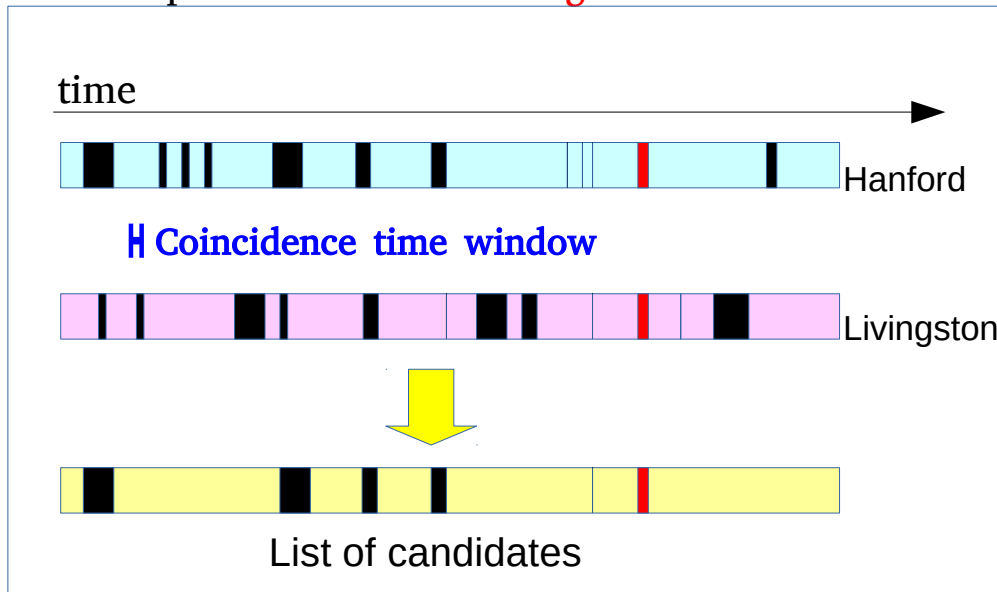
Coincidence rate:

$$R_{coinc} \sim R_H R_L \Delta t_{win}$$
$$\sim (1 \text{ Hz}) \times (1 \text{ Hz}) \times (10^{-2} \text{ s}) = 10^{-2} \text{ Hz}$$

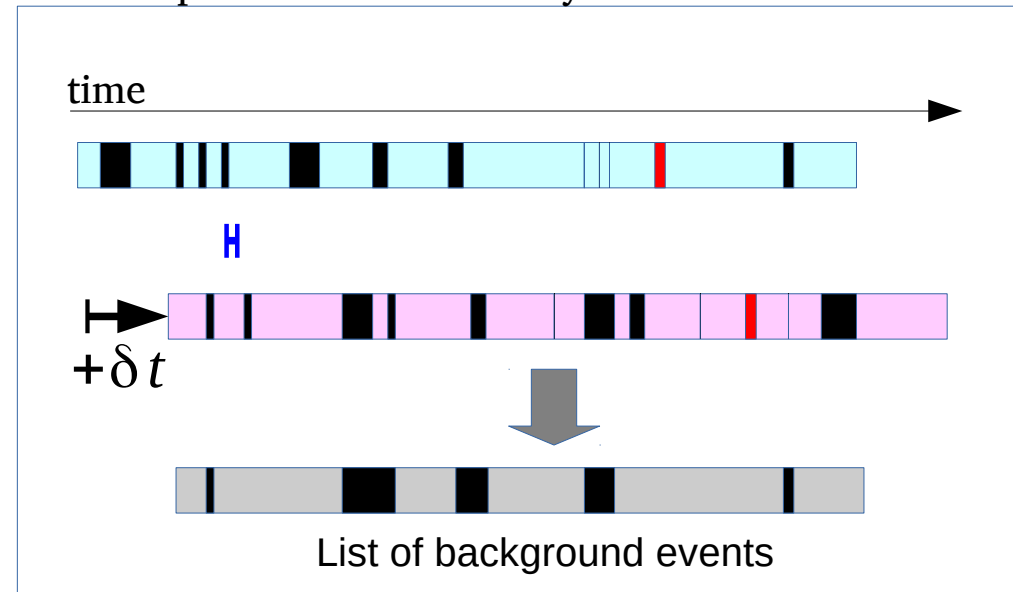
Multi-detectors coincidence

The background of a gravitational-wave search is estimated using the time-slide technique
Assumption = uncorrelated noise between detectors

True experiment = noise + **signal**



Fake experiment = noise only



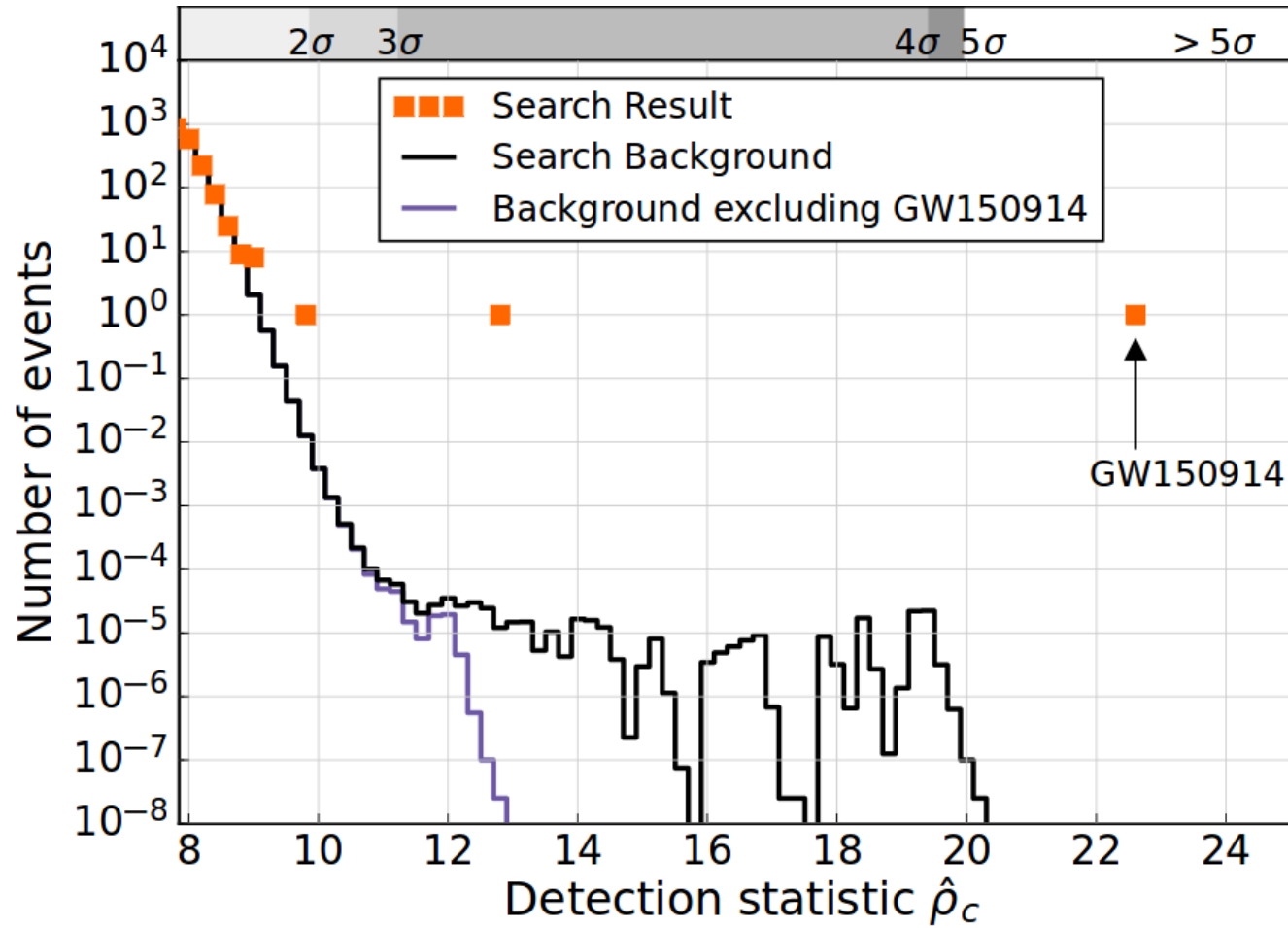
A very large number of fake experiments can be simulated using multiple offsets

LIGO O1 analysis:

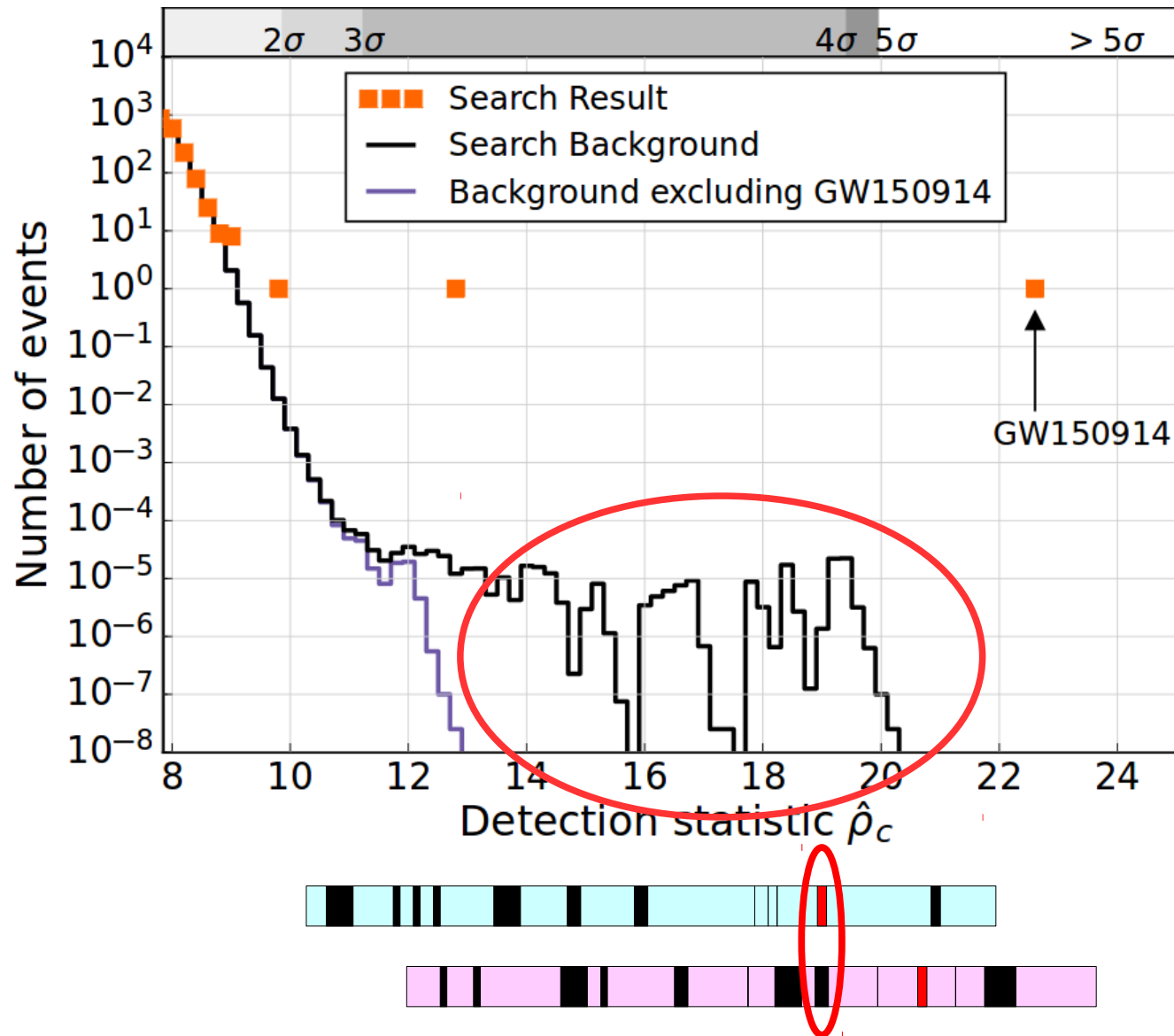
– $O(10^6)$ time offsets

→ background estimated using a fake experiment of $O(100,000)$ years

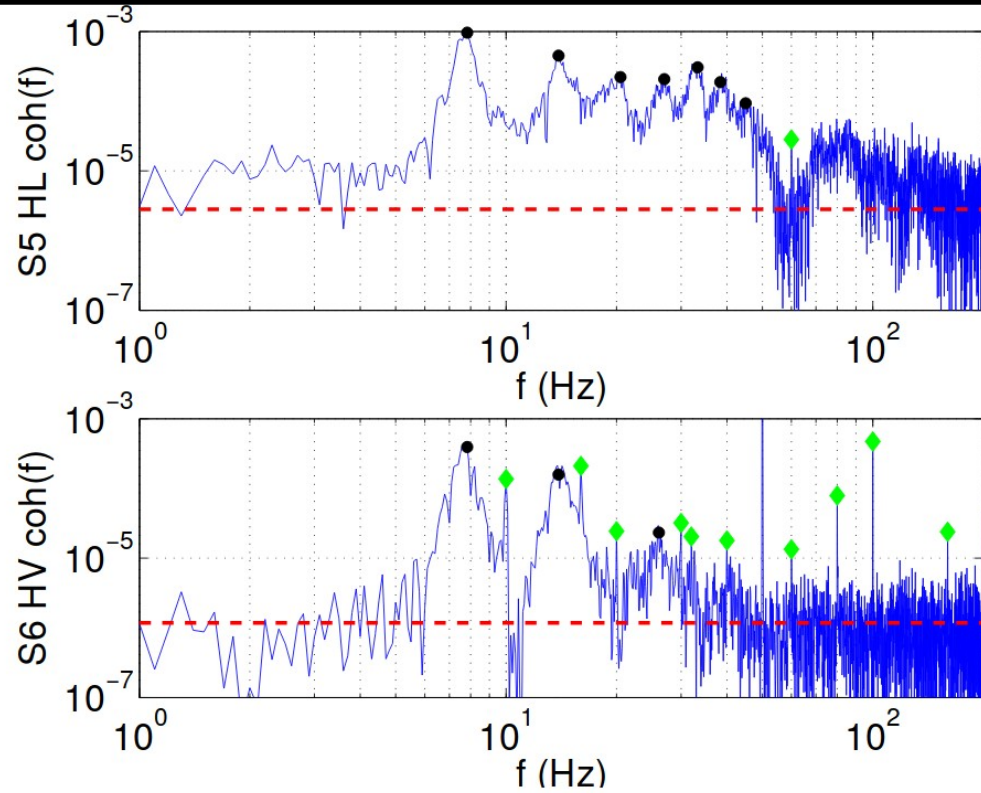
Event significance



Event significance



Correlated noise

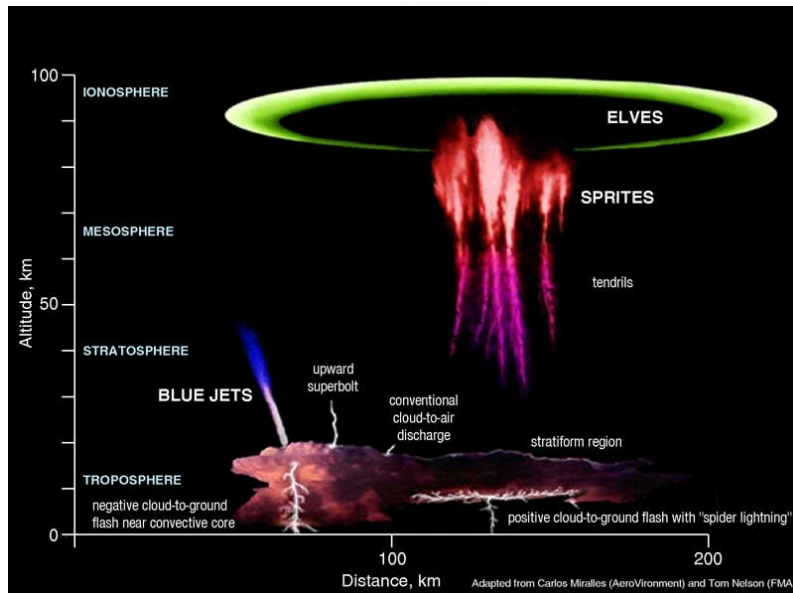
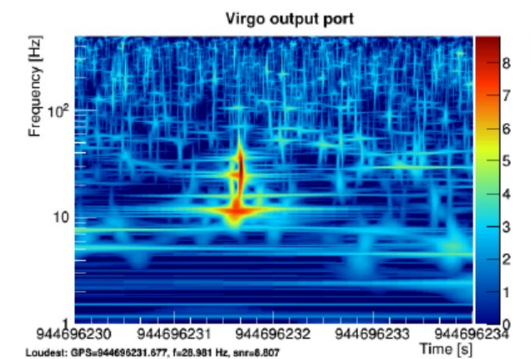
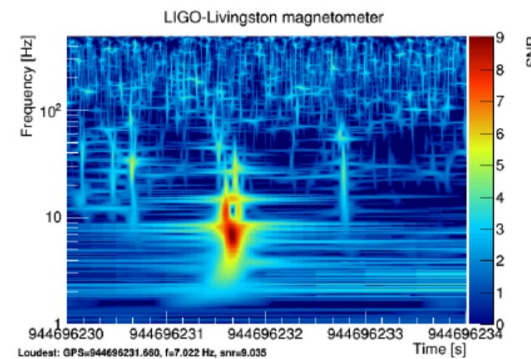
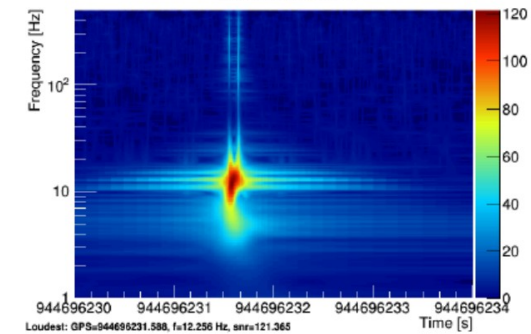
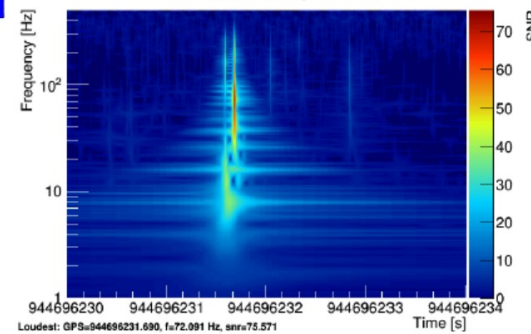


Schumann resonance



LIGO-Hanford magnetometer

Virgo magnetometer



Adapted from Carlos Miralles (AeroVironment) and Tom Nelson (FMA)

Conclusions

- First detection achieved by ground-based interferometers (LIGO-Virgo)
- A network of detectors is needed
 - to detect a gravitational-wave with confidence
 - to localize the source
 - to estimate the parameters of the source
- Analysis pipelines are used to analyze the data
- Gravitational-wave detectors are very sensitive instruments
- Multiple noise sources